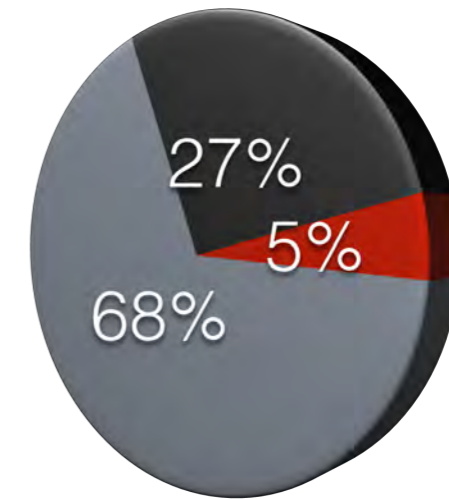
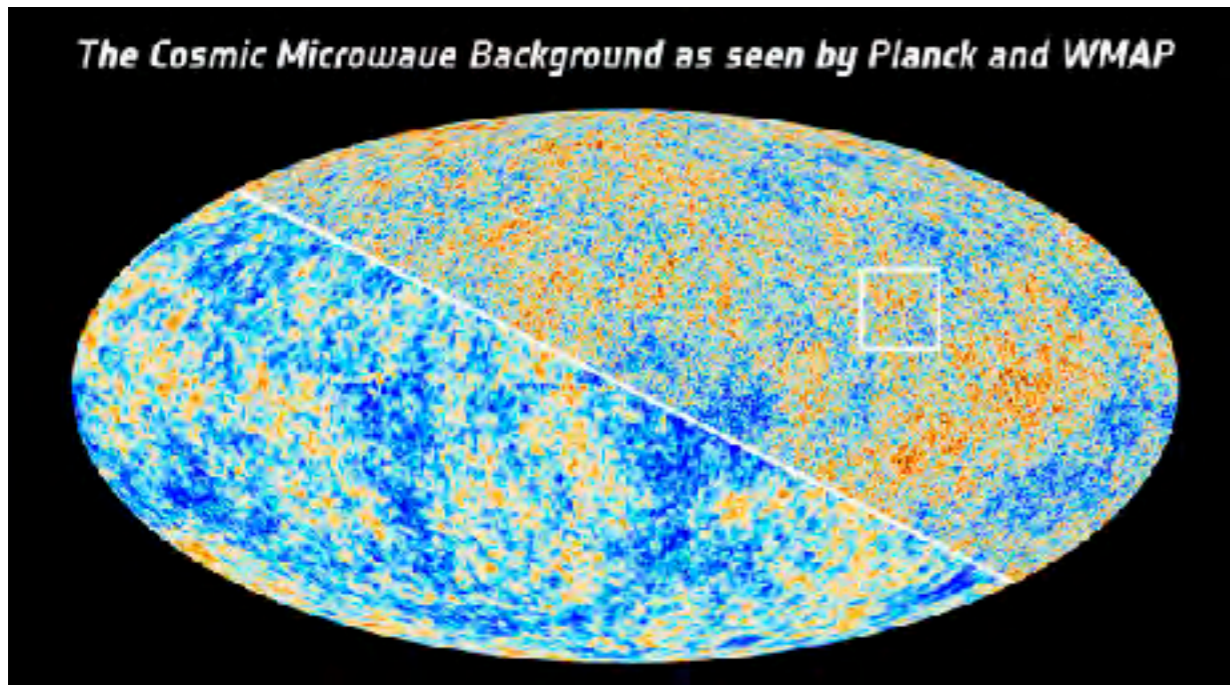


Direct detection and astrophysics

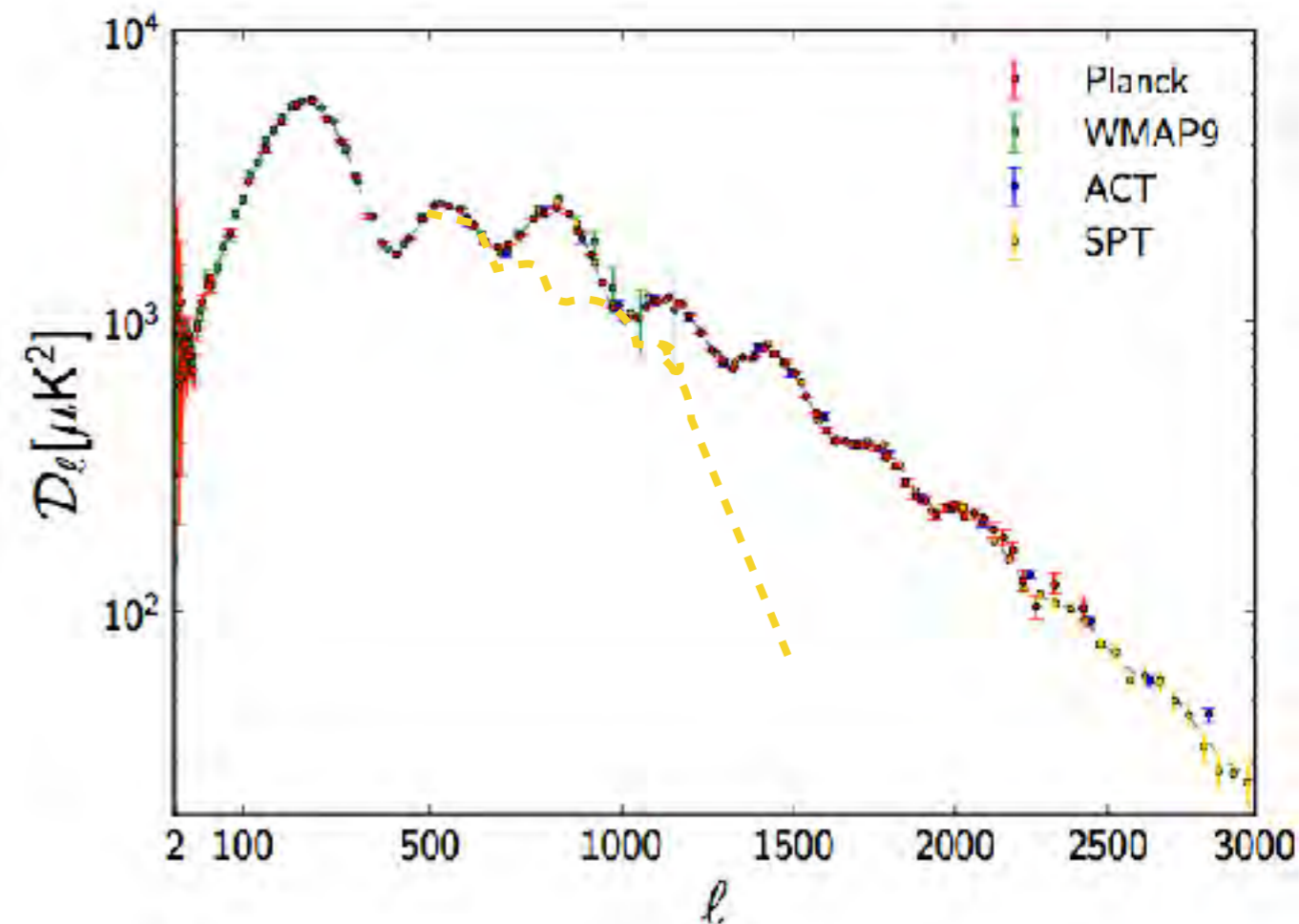
Céline Boehm

The need for non-baryonic matter



$$\begin{aligned} \dot{\theta}_b &= k^2 \psi - \mathcal{H} \theta_b + c_s^2 k^2 \delta_b - R^{-1} \dot{\kappa} (\theta_b - \theta_\gamma) \\ \dot{\theta}_\gamma &= k^2 \psi + k^2 \left(\frac{1}{4} \delta_\gamma - \sigma_\gamma \right) - \dot{\kappa} (\theta_\gamma - \theta_b), \\ \dot{\theta}_{\text{DM}} &= k^2 \psi - \mathcal{H} \theta_{\text{DM}}, \end{aligned}$$

**The CMB cannot be explained
with baryonic DM only**

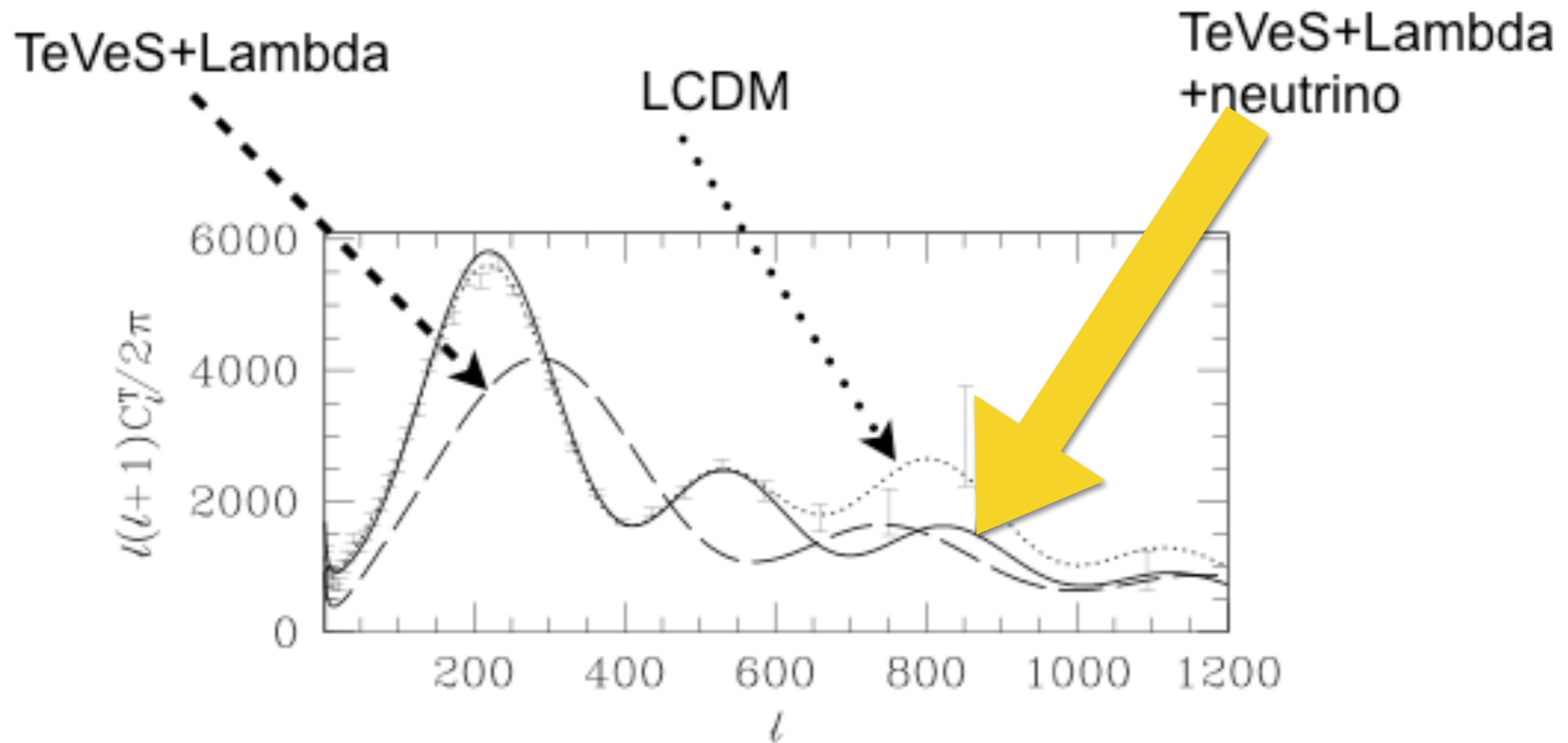


A consequence of Silk damping (Nature, 1966)

The DM particle hypothesis

Mond/Bekenstein Bekenstein [astro-ph/0403694](#)

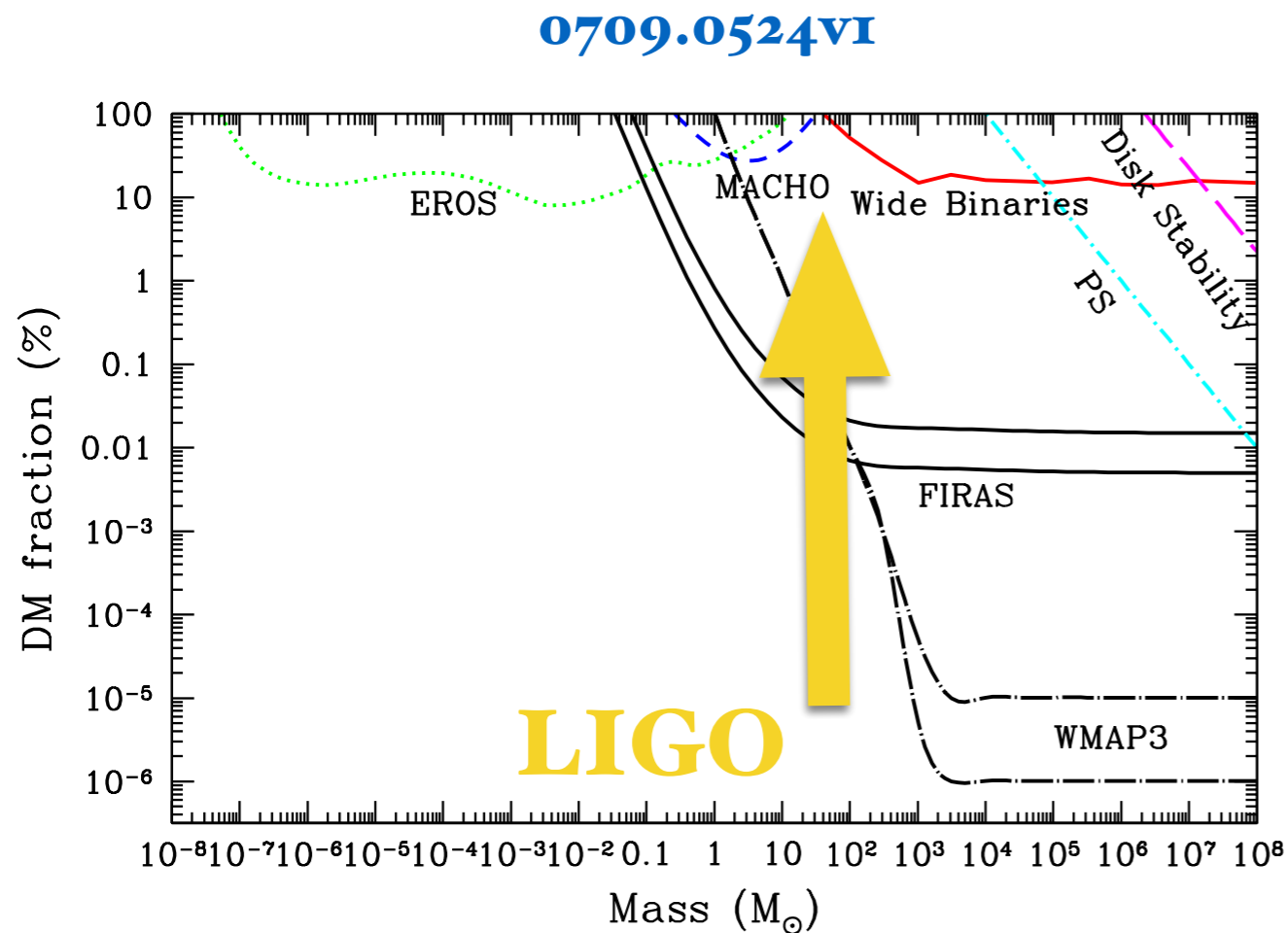
C. Skordis, D. Mota, P. Ferreira, C. Boehm : [astro-ph/0505519](#)



**Impossible to explain Planck 2015!
This is called the Silk damping (1967)**

A DM-like fluid is needed!

The need for a collisionless fluid



Primordial Black Holes ???

[arXiv:1603.00464](https://arxiv.org/abs/1603.00464)

[arXiv:1501.07565](https://arxiv.org/abs/1501.07565)

[arXiv:1607.06077](https://arxiv.org/abs/1607.06077)

Ways to evade CMB limits

[arXiv:1612.05644](https://arxiv.org/abs/1612.05644)

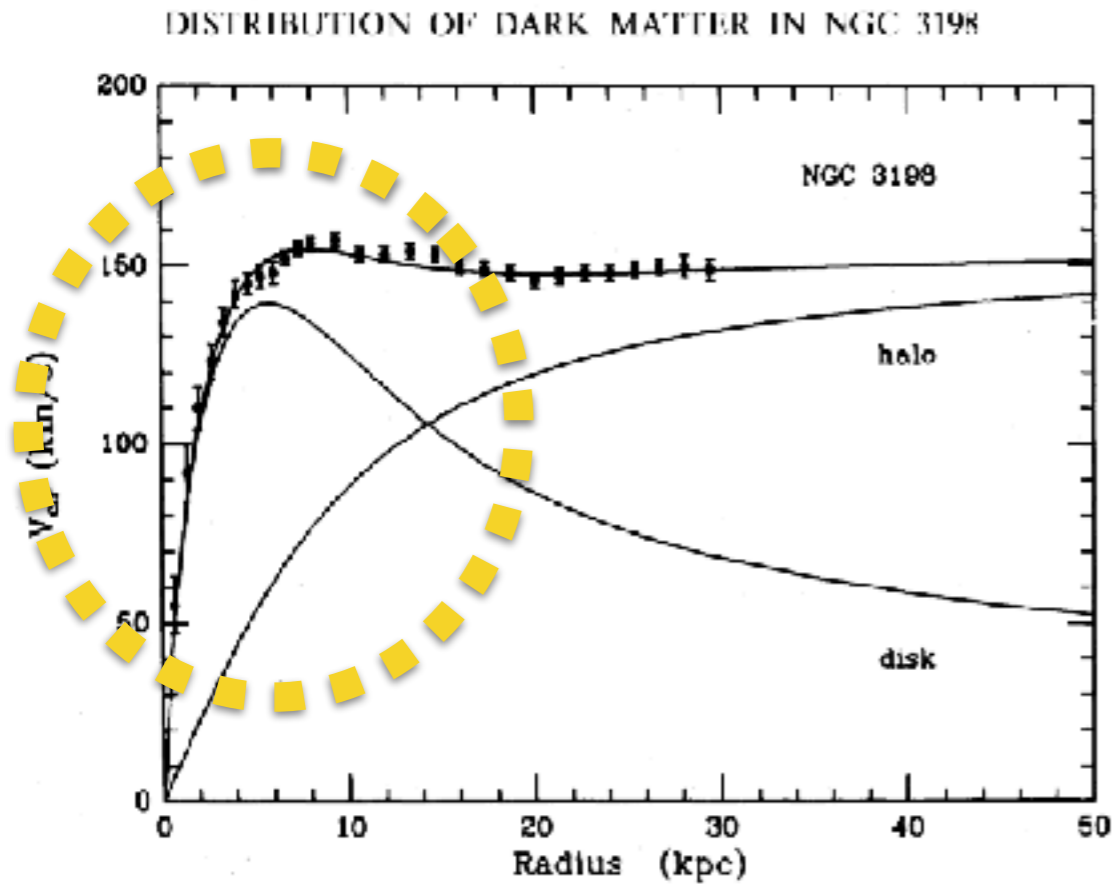
> 100 M_{sol} ruled out as main DM component

**But we still need some sort of dark matter
(at least ~ a collisionless fluid)**

WIMPs are still the simplest explanation

How do we test the particle hypothesis?

The DM halo



$$v_c^2 = \frac{G M(r)}{r}$$

$$M(r) = \int 4\pi^2 \rho(r) dr^3$$

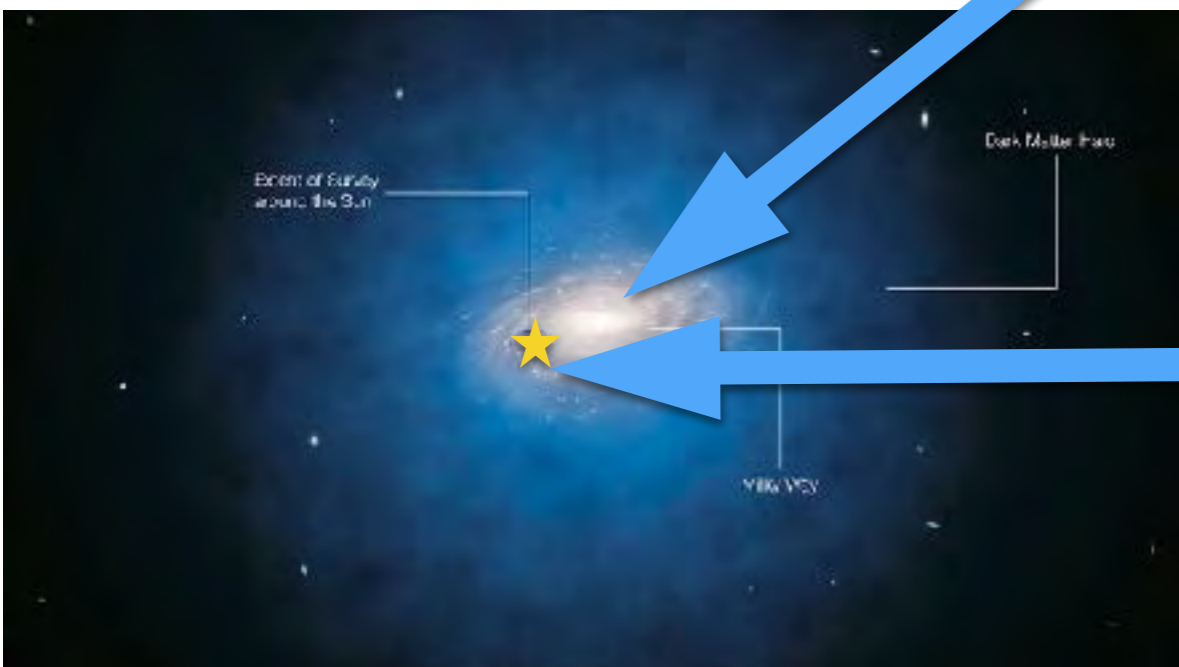
$$M(r) = M_{dm} + M_{baryons}$$

DM is a subcomponent
Its density is the highest

great for Indirect Detection

DM is dominant.
Its density is low

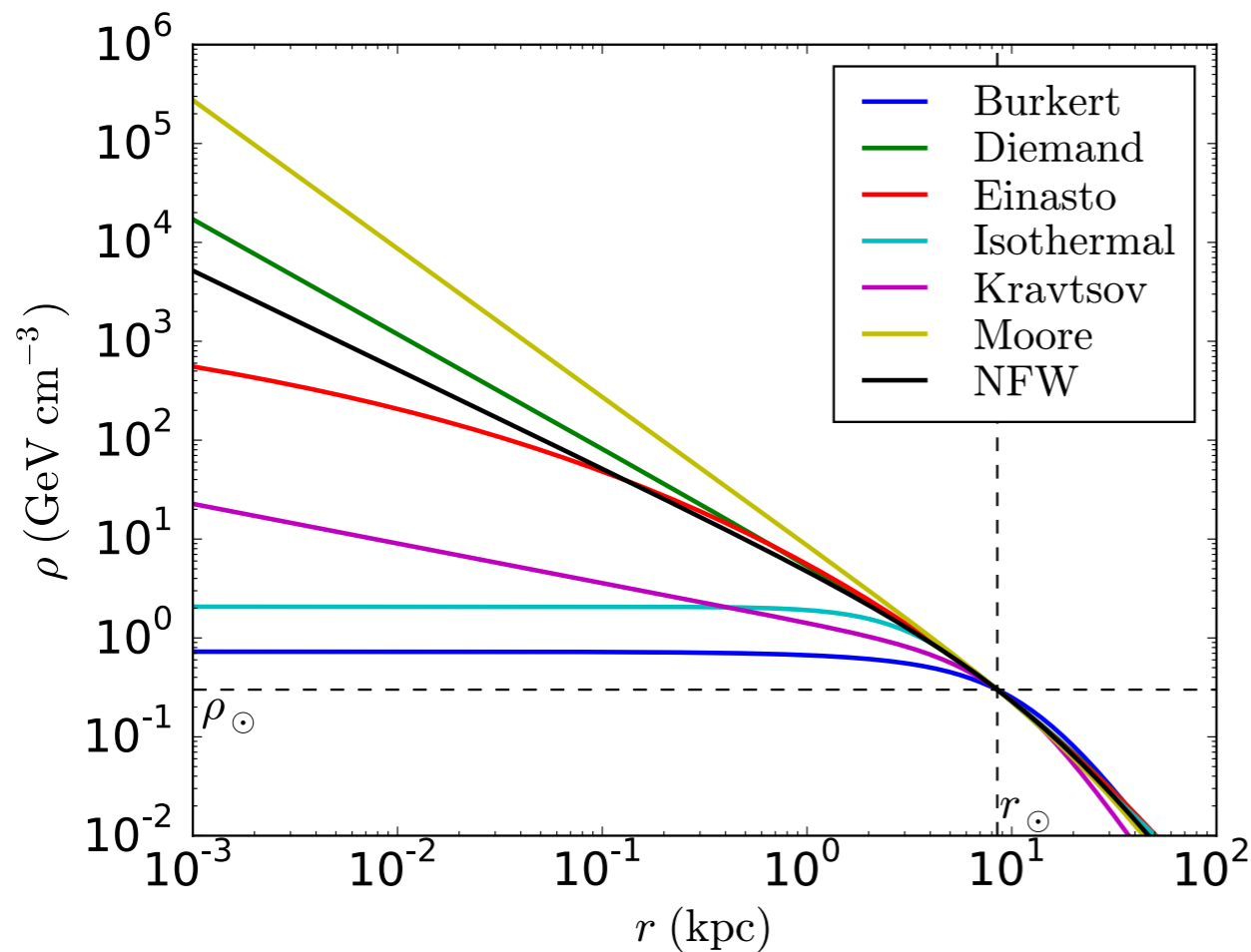
great for direct Detection



Potential issues

Indirect Detection

$$\phi \propto \int dl \rho(r)^{1,2}$$



Courtesy: T. Lacroix

$$\rho_{\text{NFW}_{\text{gen}}}(r) = \rho_0 \left(\frac{r}{r_0} \right)^{-\gamma} \left[1 + \left(\frac{r}{r_0} \right)^\alpha \right]^{-\frac{\beta-\gamma}{\alpha}}$$

Direct Detection

$$\frac{dR}{dE} = \frac{\sigma(q)}{2m\mu^2} \rho \eta(E, t)$$

$$\eta(E, t) = \int_{v_{\min}(E)}^{\infty} \frac{f(v, u_e(t))}{v} d^3v$$

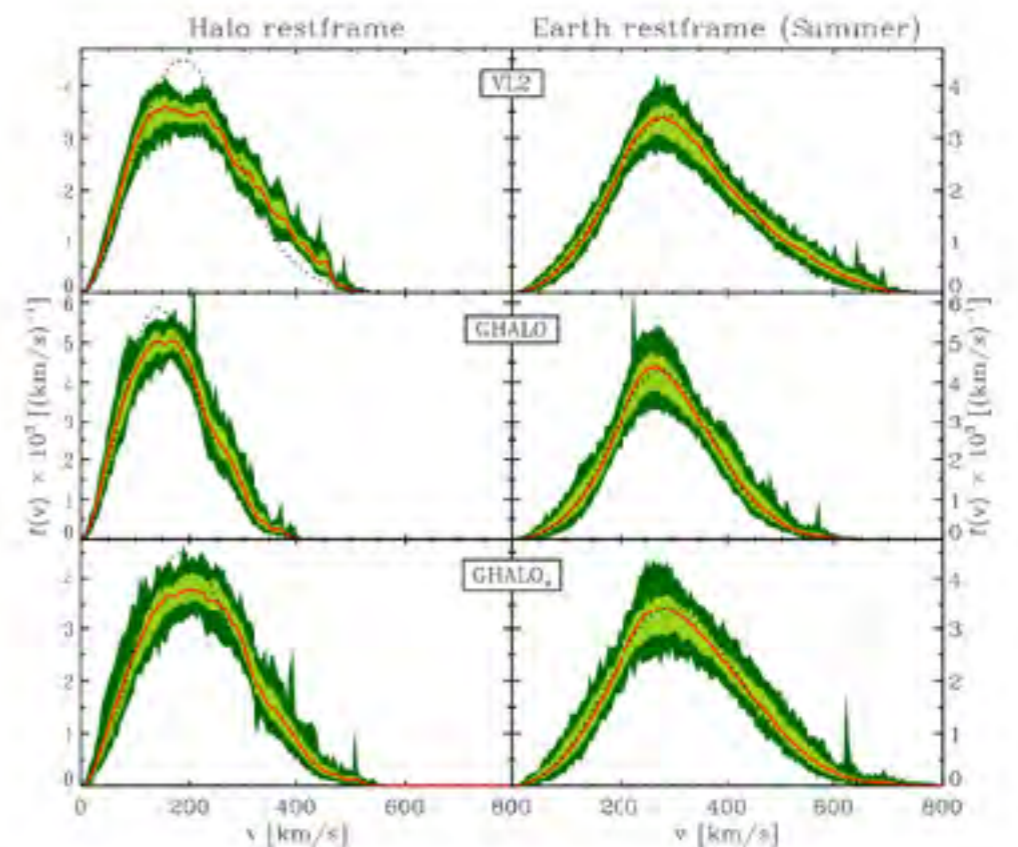
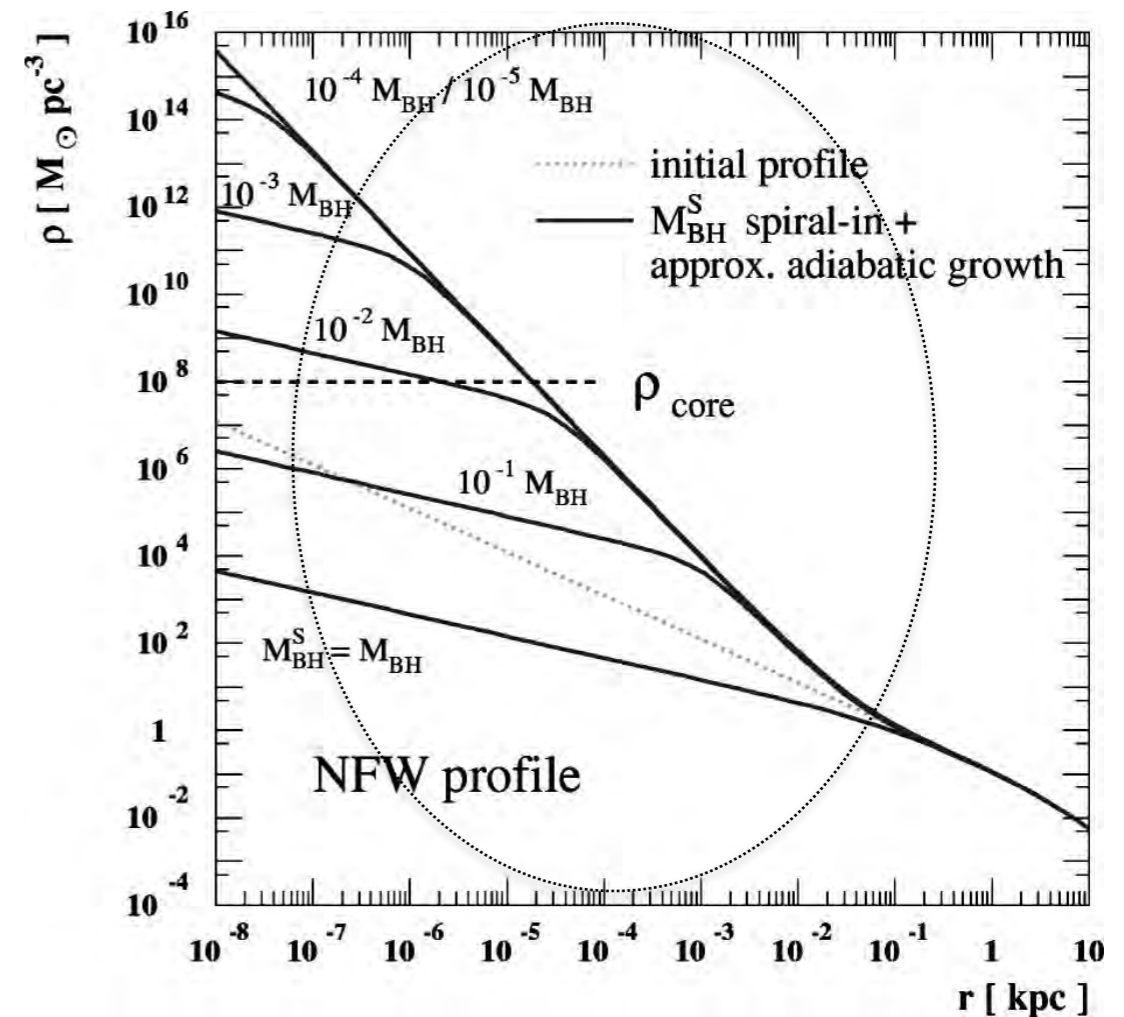
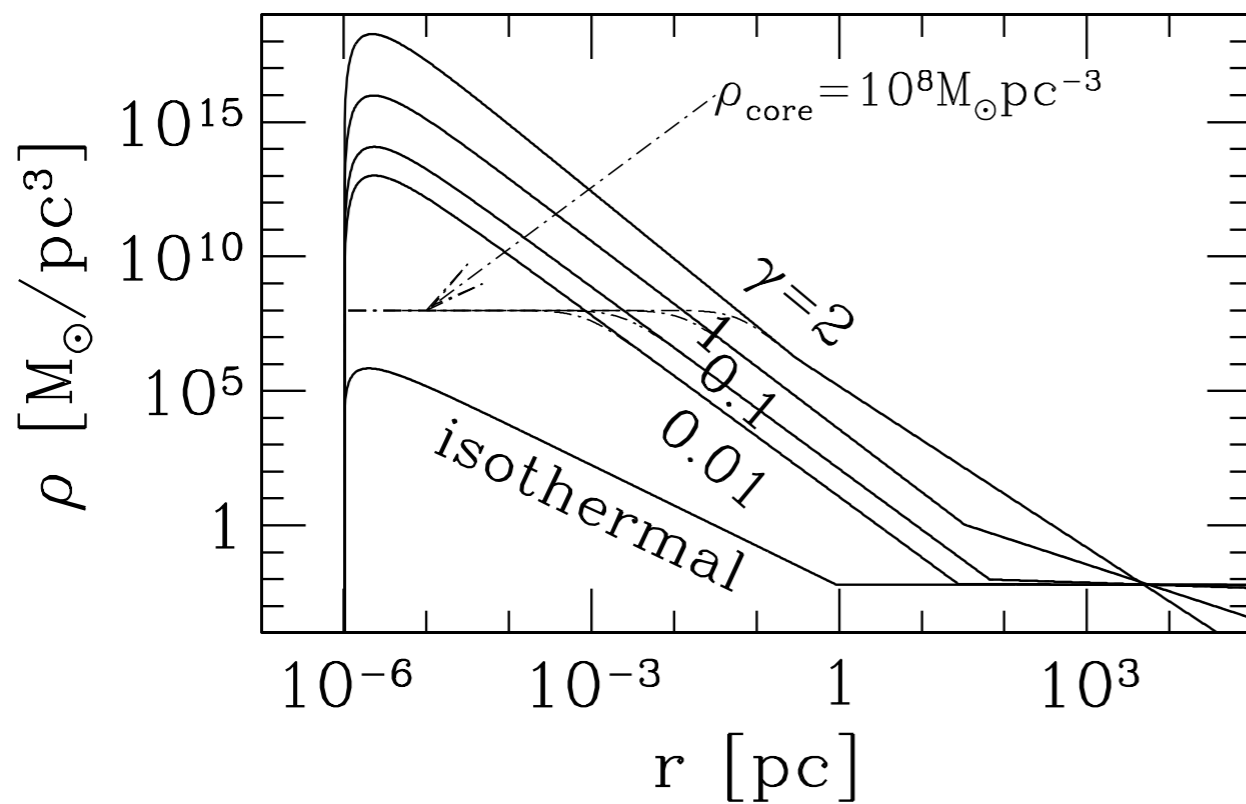


Figure 2. Velocity distribution functions: the left panels are in the host halo's restframe, the right panels in the restframe of the Earth on June 2nd, the peak of the Earth's velocity relative to Galactic DM halo. The solid red line is the distribution for all particles in a 1 kpc wide shell centered at 8.5 kpc, the light and dark green shaded regions denote the 68% scatter around the median and the minimum and maximum values over the 100 sample spheres, and the dotted line represents the best-fitting Maxwell-Boltzmann distribution.

see A. Peter's talk

DM accretion near Black Holes



Condition of formation of spikes:

- * BHs at the center of galaxies can grow adiabatically
- * Adiabatic growth inside a population of stars enhances the density of stars

Why not enhancement of DM density?

Ipsier & Sikivie (1987): isothermal $\rightarrow r^{-3/2}$, Gondolo & Silk (1999) : NFW $\rightarrow 7/3$

Electromagnetic emission from cosmic rays from DM

Cluster of galaxies

Milky Way

Dwarf galaxies

AGNs (Cen A, ESO)



cosmic rays

gamma rays

radio/submm

decay

$DM \rightarrow SM SM (SM)$

annihilations

$DM DM \rightarrow SM SM$

Indirect detection signatures

Injection of cosmic rays at high energy (> keV!)

If $m_{\text{DM}} \gg \text{MeV}$

Prompt emission (gamma-ray)

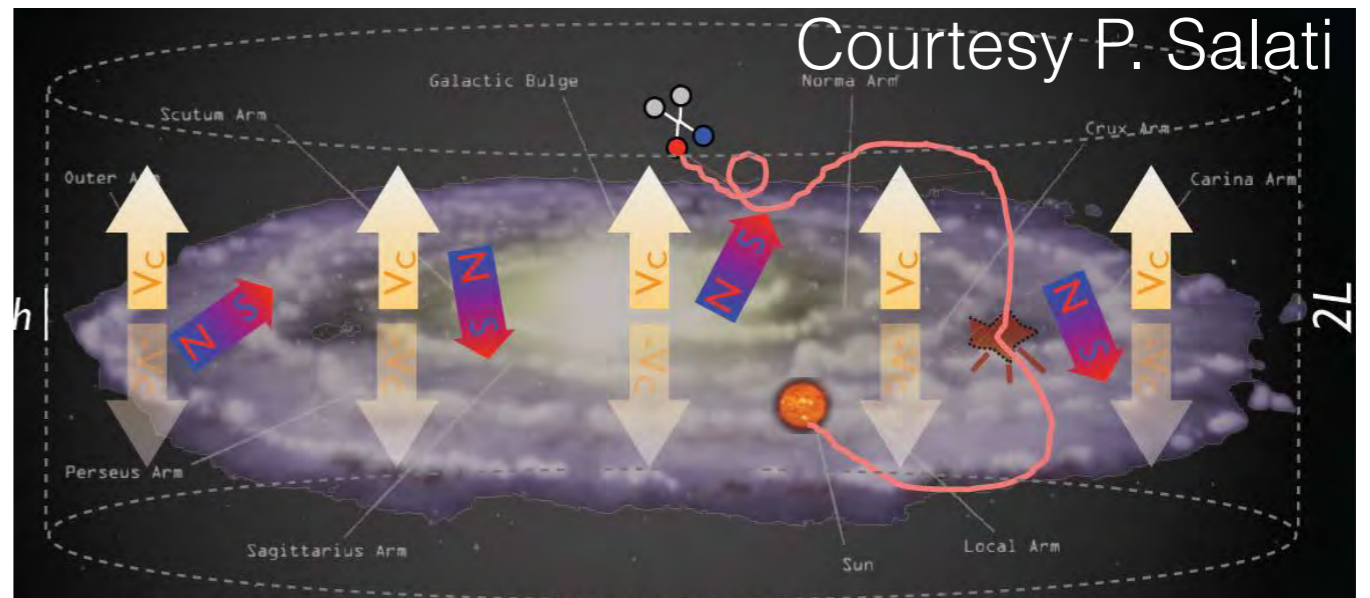
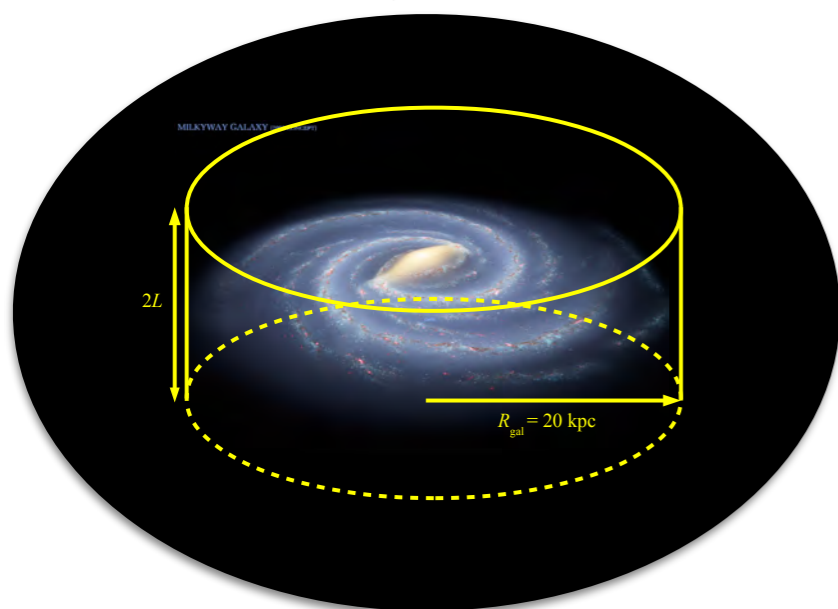
$$\frac{d\phi}{dE} = \frac{1}{8\pi} \left(\frac{\sigma v}{m_{\chi^2}^2} \right) \sum_i \text{BR}_i \frac{dN_i}{dE} \xi^2 \int dl \rho_{\chi}^2(l) \Rightarrow \Phi_{\text{prompt}} \propto \frac{\sigma v}{m_{\text{DM}}^2} \int dl \rho^2(l)$$

Only depends on the DM profile

Inverse Compton (gamma-ray)
Synchrotron (radio/submm)

- 1) Spatial diffusion
- 2) Energy losses

“propagation” of cosmic rays
gives rise to typical energy spectra



Submm constraints

arXiv:1105.4689

our Milky Way

Astrophysical
sources

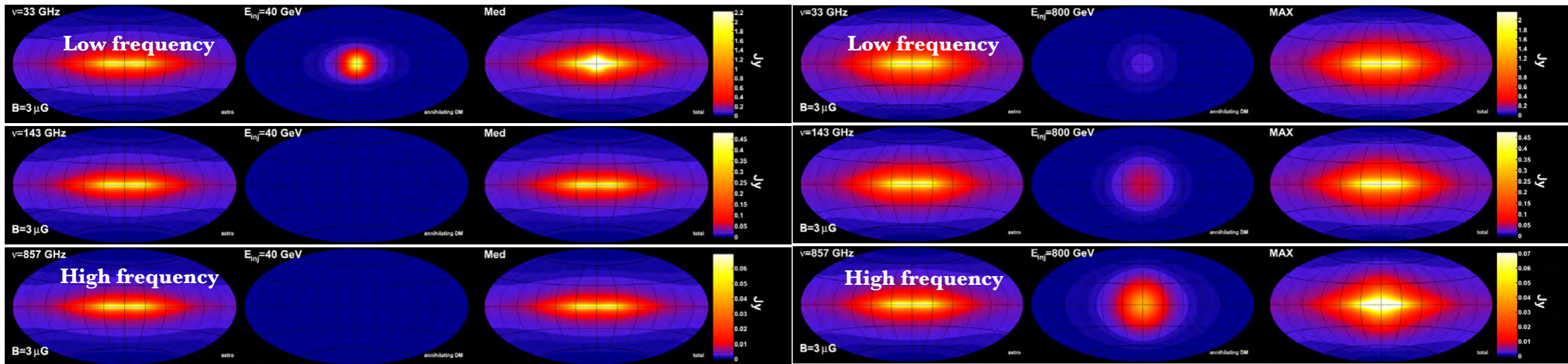
DM

Sum

Astrophysical
sources

DM

Sum



40 GeV DM

Expected signature in LFI

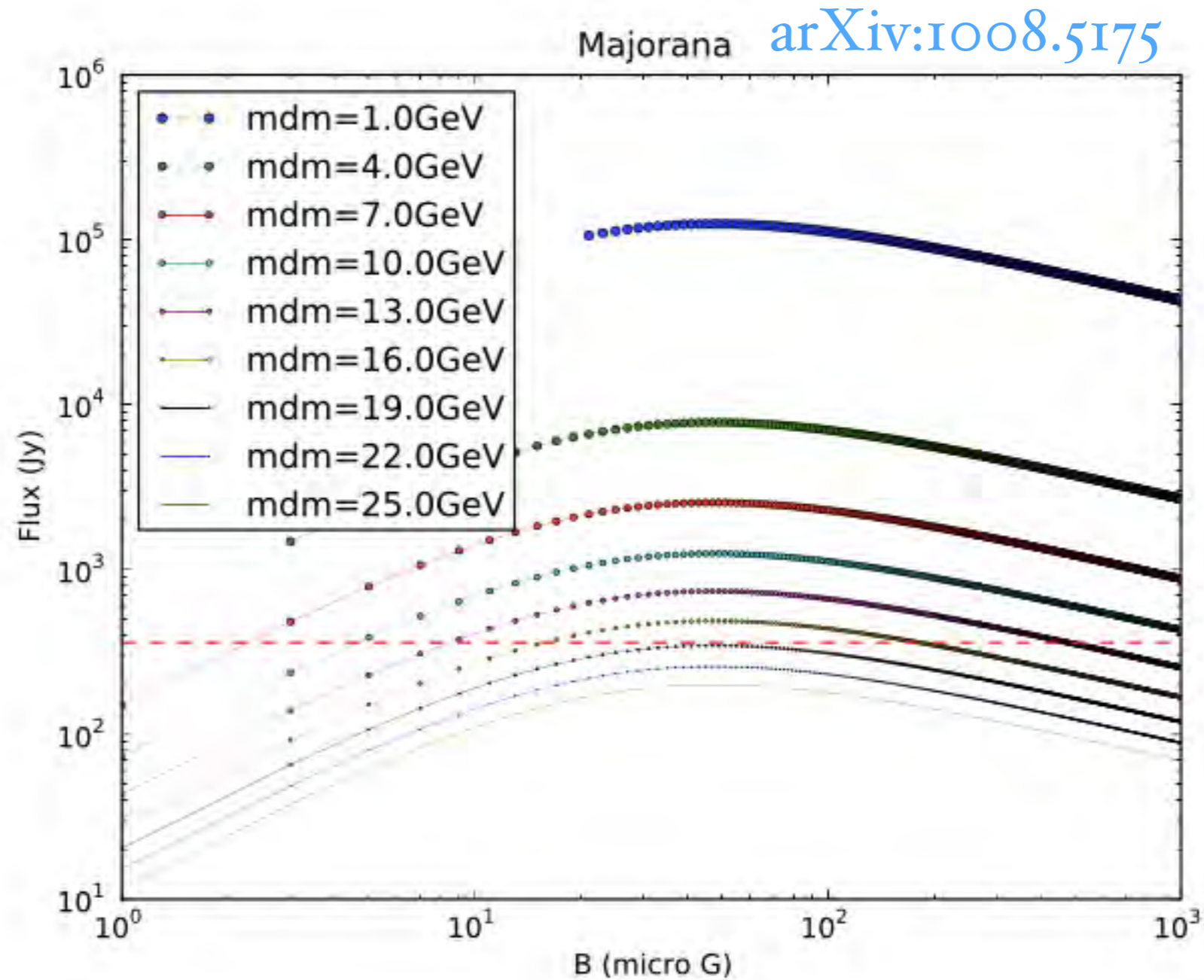
800 GeV DM

Expected signature in HFI

No signal!

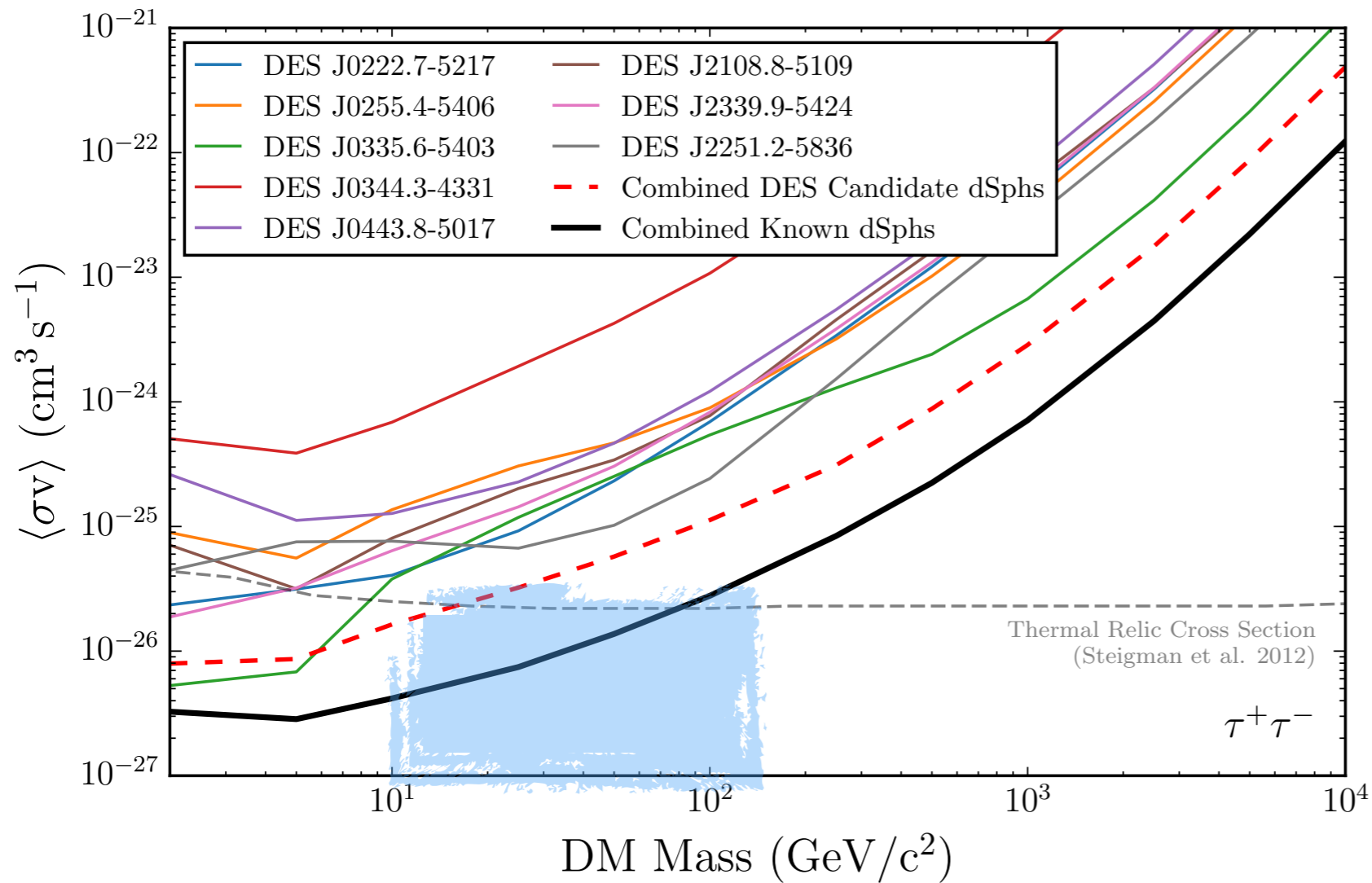
Radio constraints our Milky Way

330 MHz



Excludes up to 10 GeV particles for normal B field values

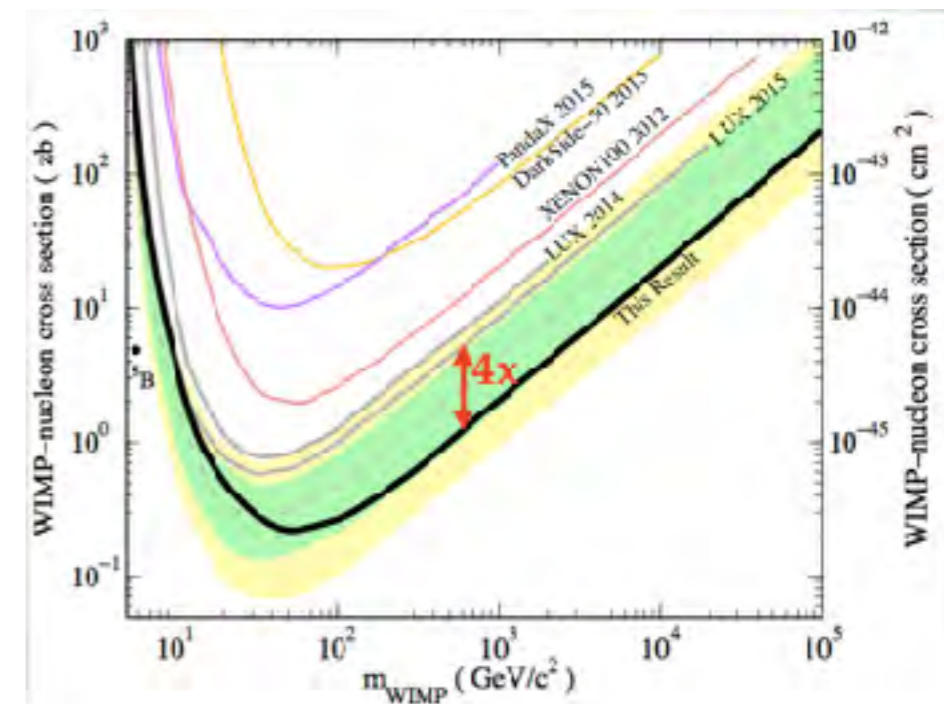
Gamma-rays in dSphs



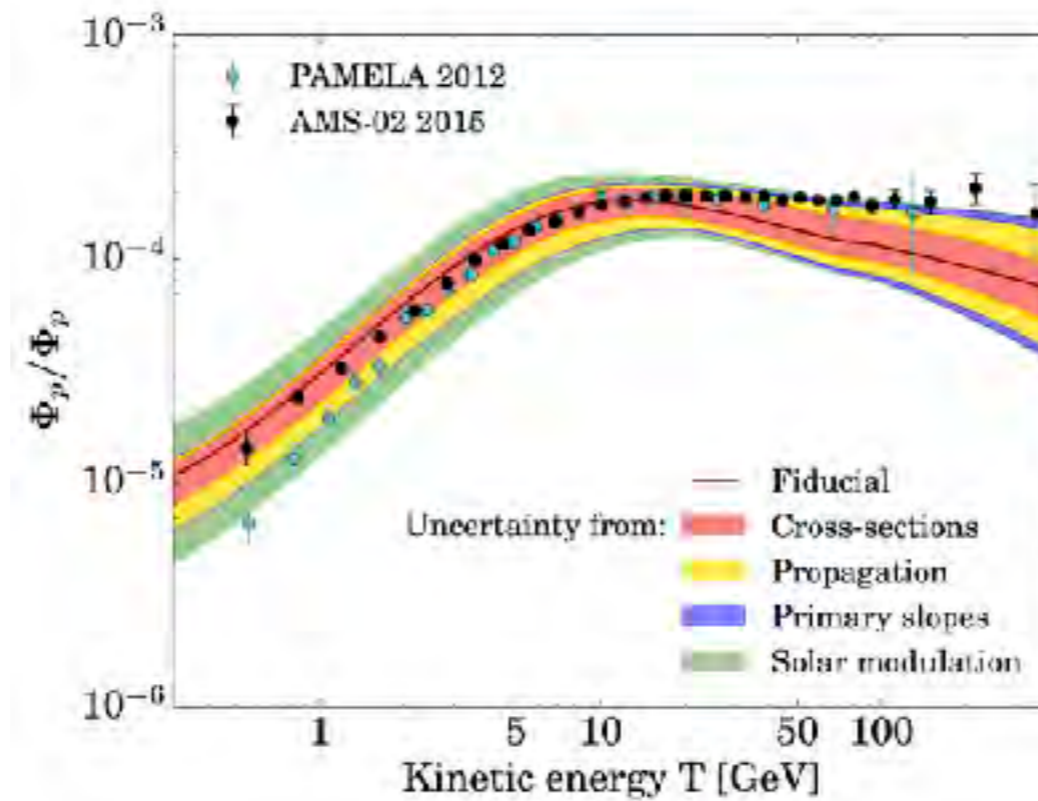
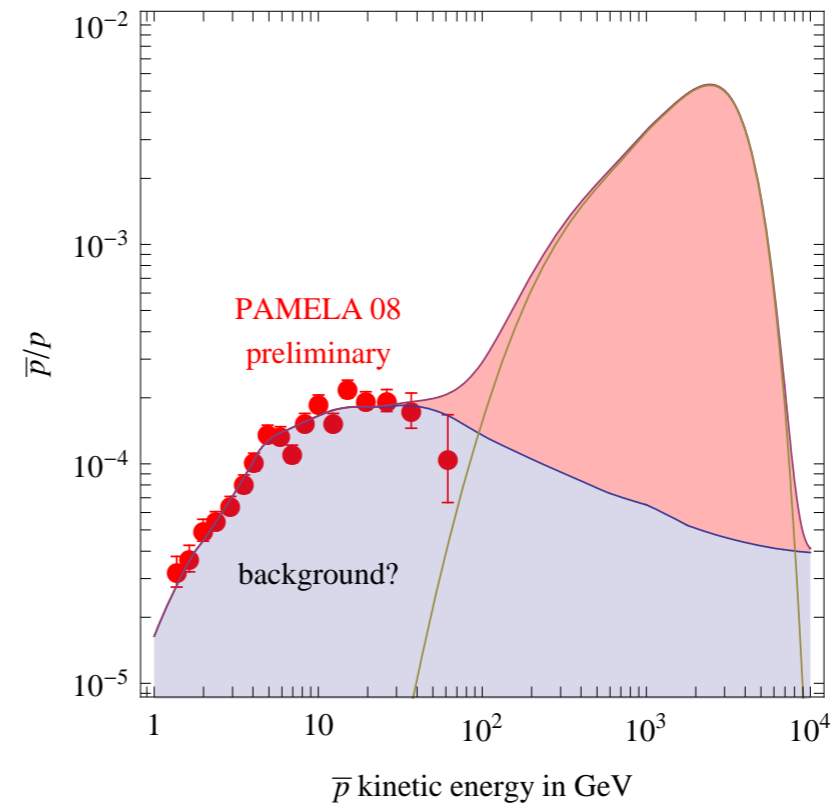
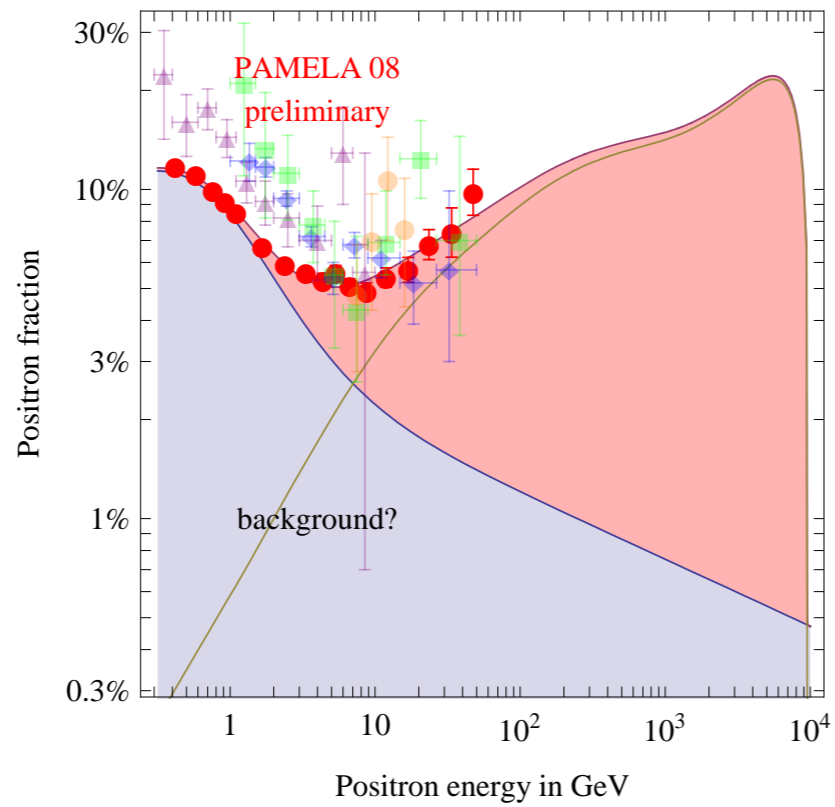
Thermal DM already partially excluded!

~ 3 order of magnitude to catch up DD limits though the process aren't the same!

I'll comment again on this!

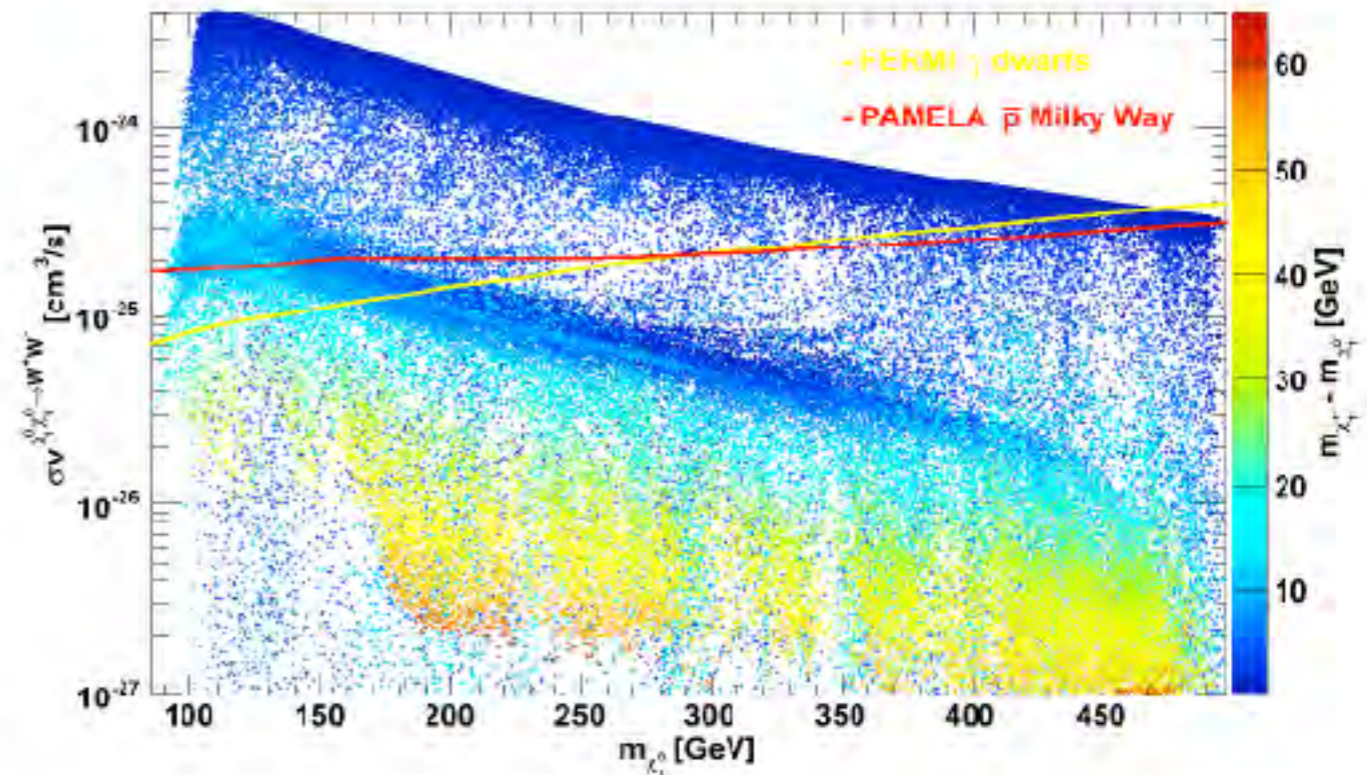
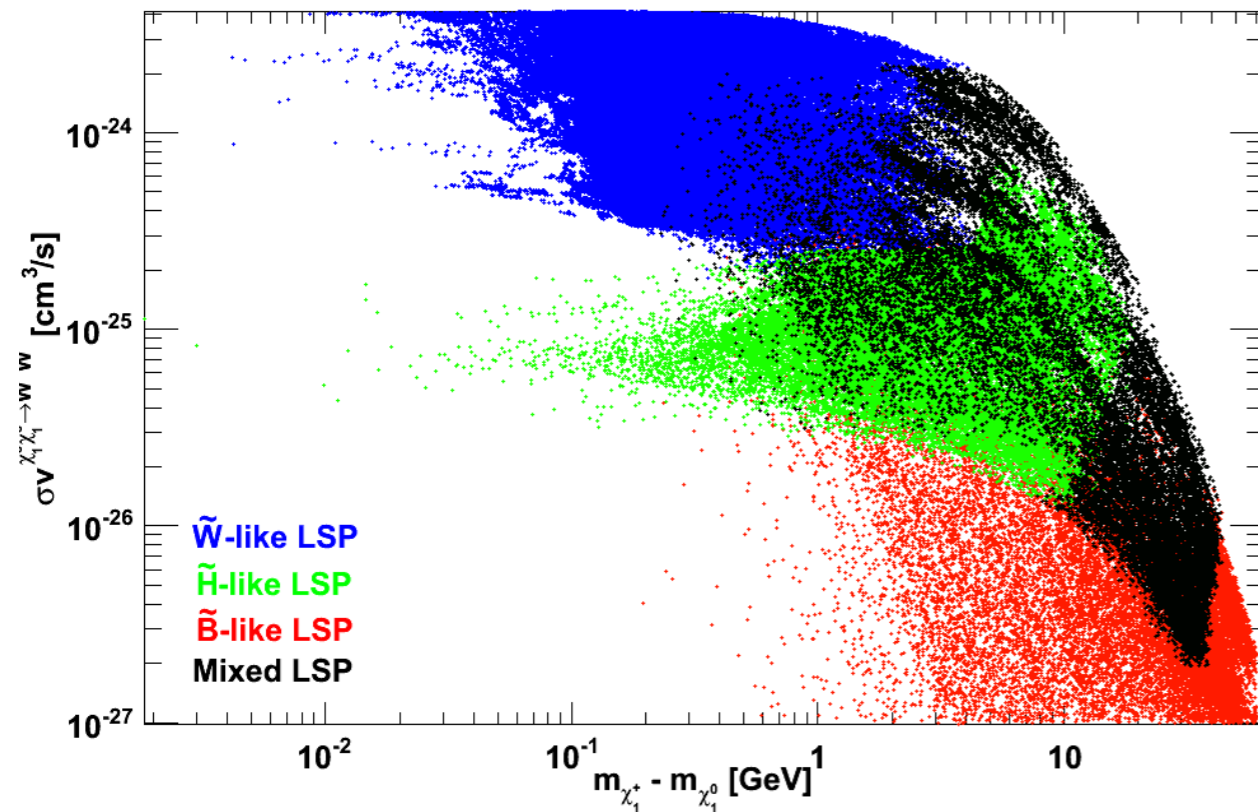


Antiprotons



Constraints on neutralino annihilating into $W+W-$

arXiv:1208.5009



We can exclude certain cross section and masses

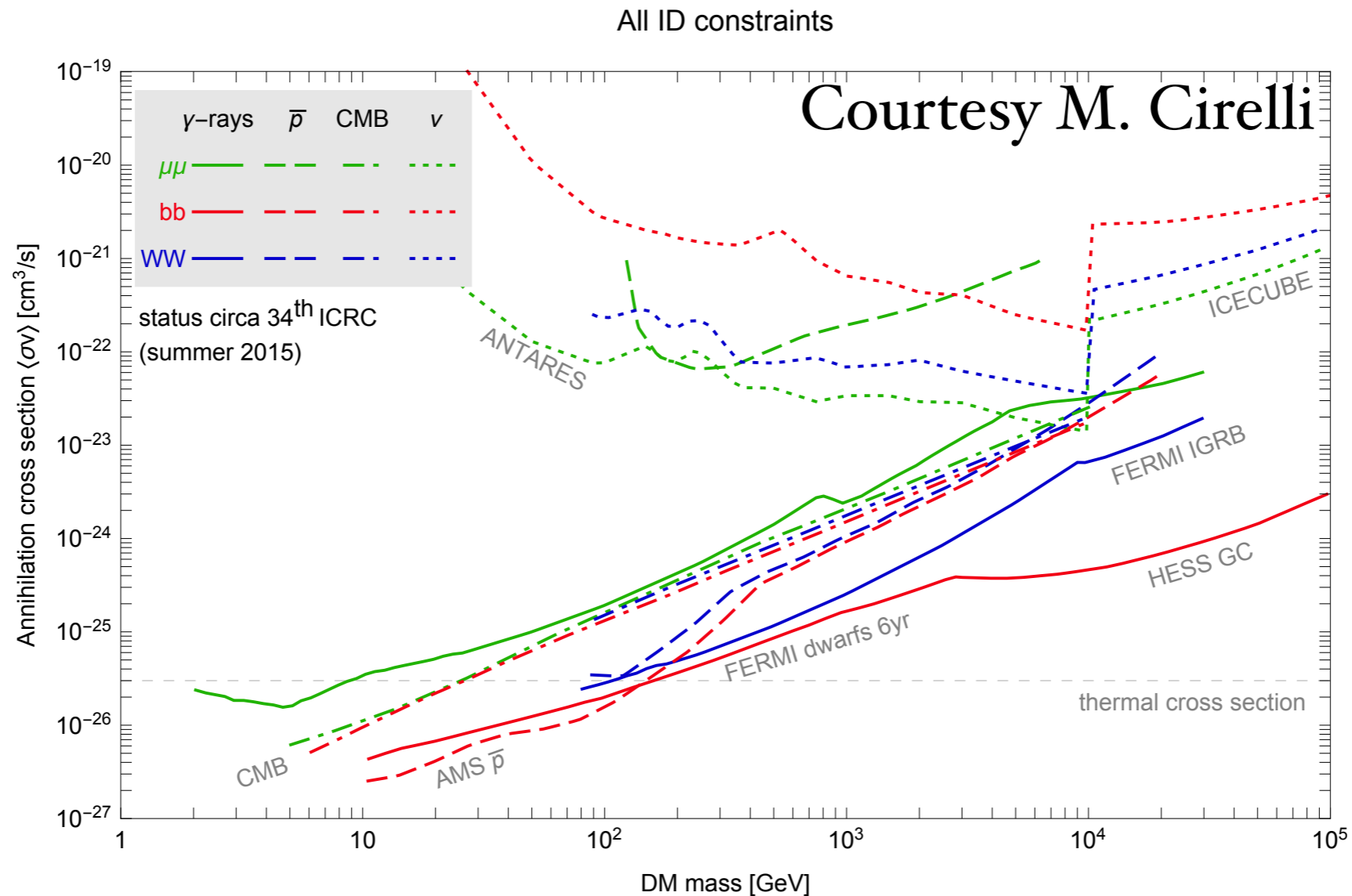
- ✓ One can exclude values of the annihilation cross section versus the neutralino mass
- ✓ From the value of the annihilation cross section, one can exclude the neutralino composition

One can therefore exclude certain neutralino composition on the sole basis of the anti-proton flux predicted in these DM scenarios.

Also recently [arXiv:1401.6212](https://arxiv.org/abs/1401.6212)

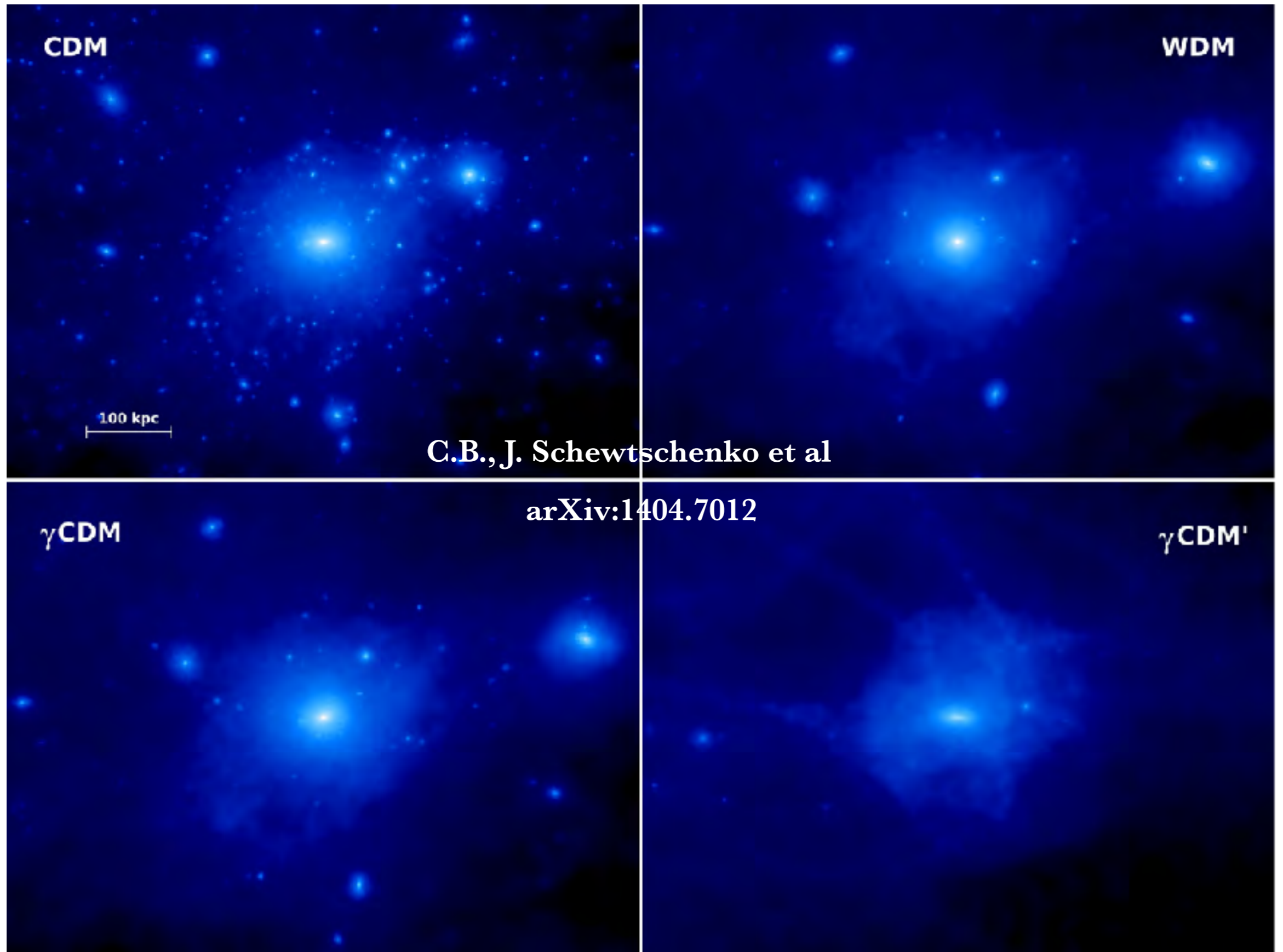
Gamma-rays

dSphs + MW + CMB



In reality one can constrain light DM (< 10 GeV) too!

$$\sigma v < 10^{-31} \left(\frac{m_{\text{DM}}}{\text{MeV}} \right)^2 \text{ cm}^3/\text{s} \quad \text{for } < \mathbf{O(\text{GeV})}$$



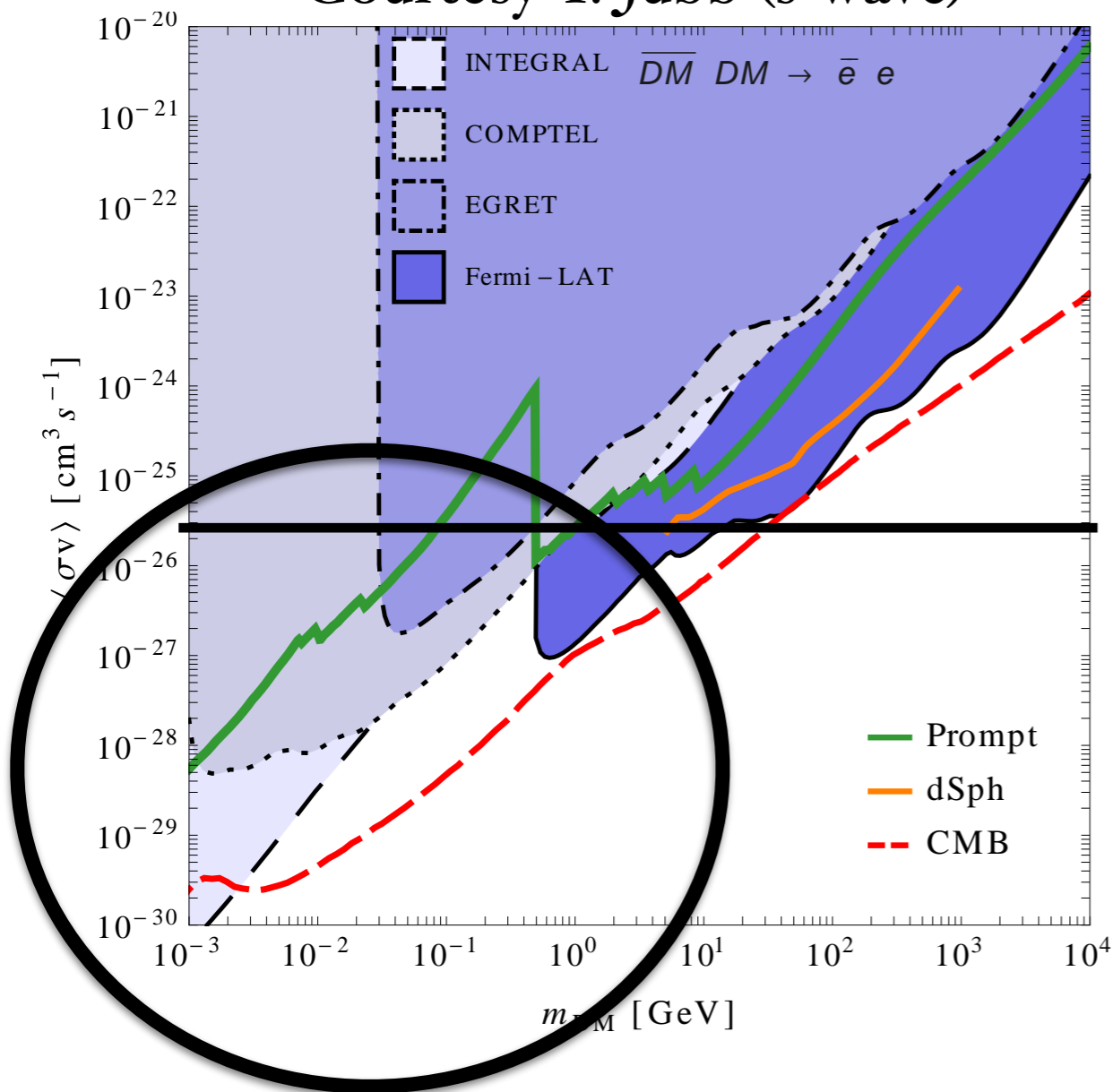
$$\sigma_{\text{DM}-\gamma} \lesssim 10^{-33} \left(\frac{m_{\text{DM}}}{\text{GeV}} \right) \text{cm}^2$$

same with neutrinos

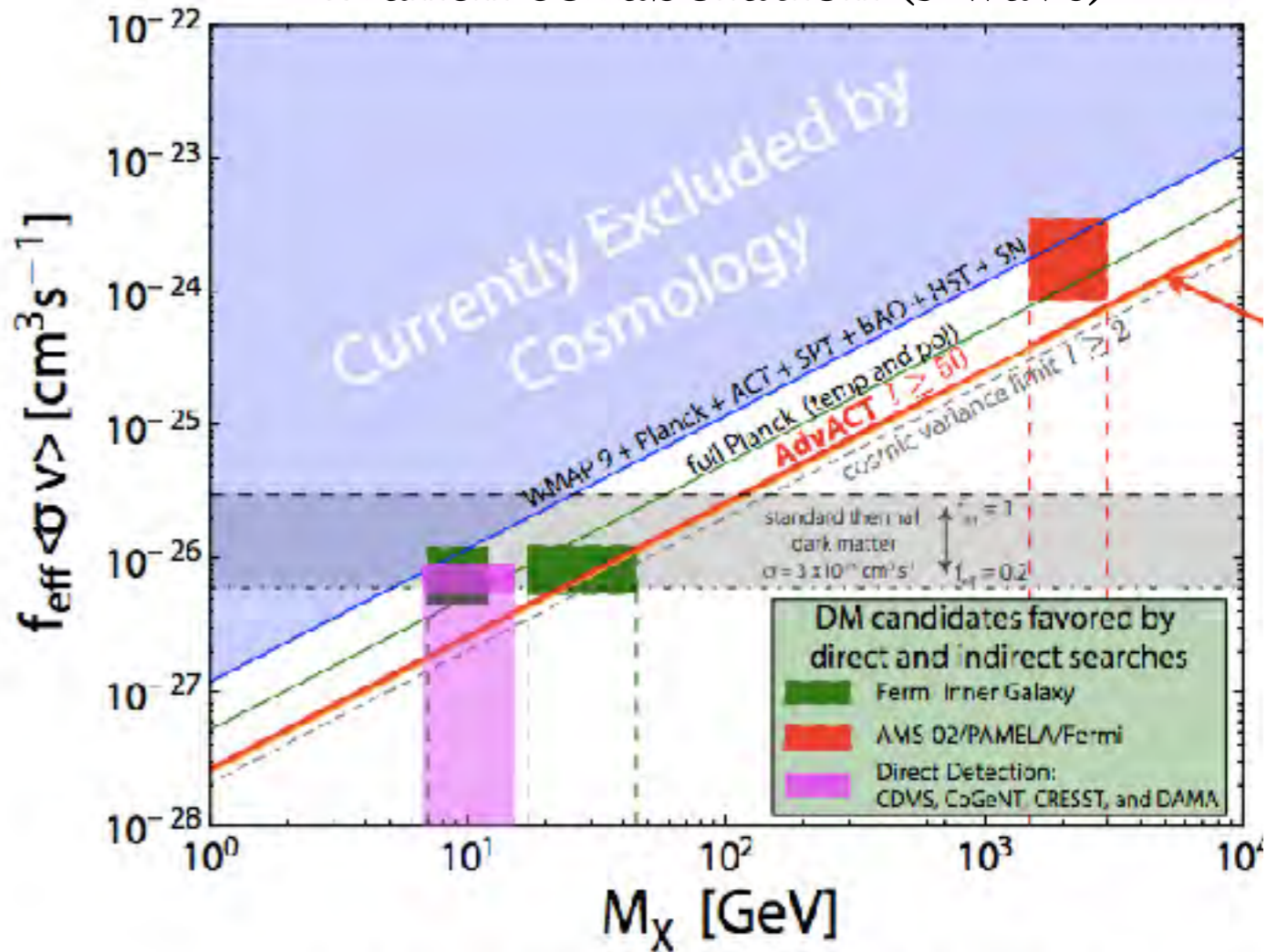
Indirect detection limits have strong impact on $m_{\text{DM}} < 10 \text{ GeV}$

DM can be lighter than a few GeV! But ...

Courtesy T. Jubb (s-wave)



Planck collaboration (s-wave)

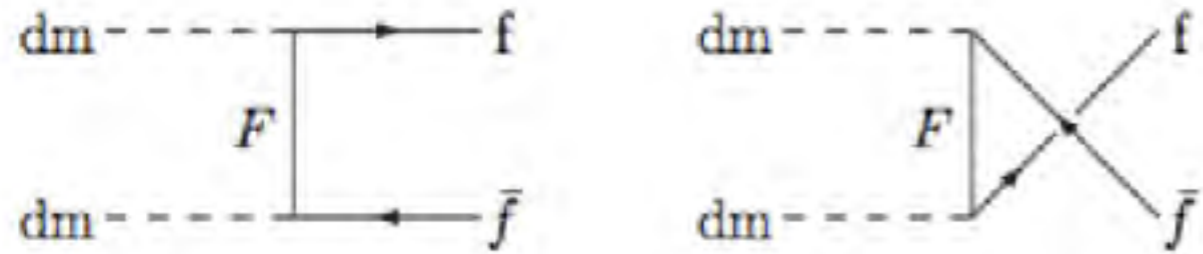


Madhavacheril, NS, Slatyer 2014, PRD, (1310.3815)

For light DM, the annihilation cross section into electrons needs to be very suppressed!

Indirect detection limits have strong impact on $m_{\text{DM}} < 10 \text{ GeV}$

The cross section can be independent of the DM mass!



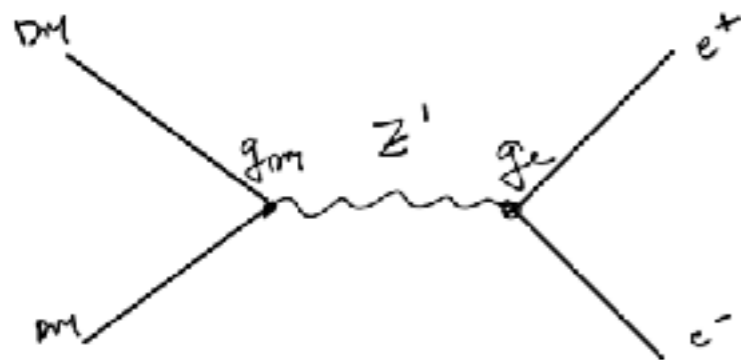
non chiral couplings

$$\sigma v \propto \frac{1}{m_F^4} \left((C_l^2 + C_r^2) m_f + 2C_l C_r m_F \right)^2$$

CB, P. Fayet, hep-ph/0305261

Feng & Kumar (0803.4196)

The mediator can be very light!



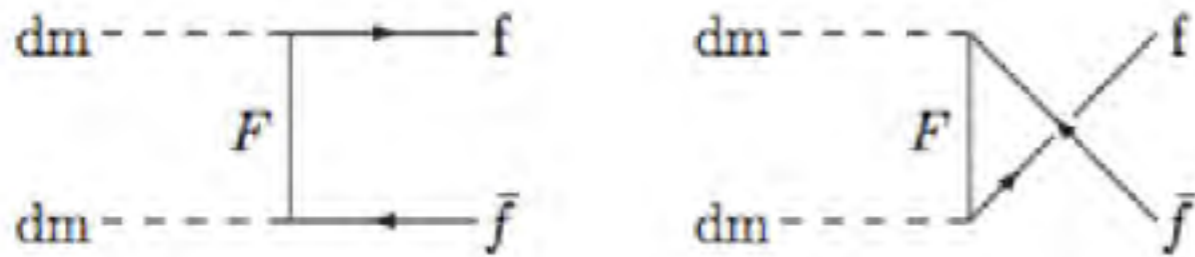
$$\sigma v \propto v^2 \frac{m_{\text{DM}}^2}{m_{Z'}^4} g_{\text{DM}}^2 g_e^2$$

CB, P. Fayet, hep-ph/0305261

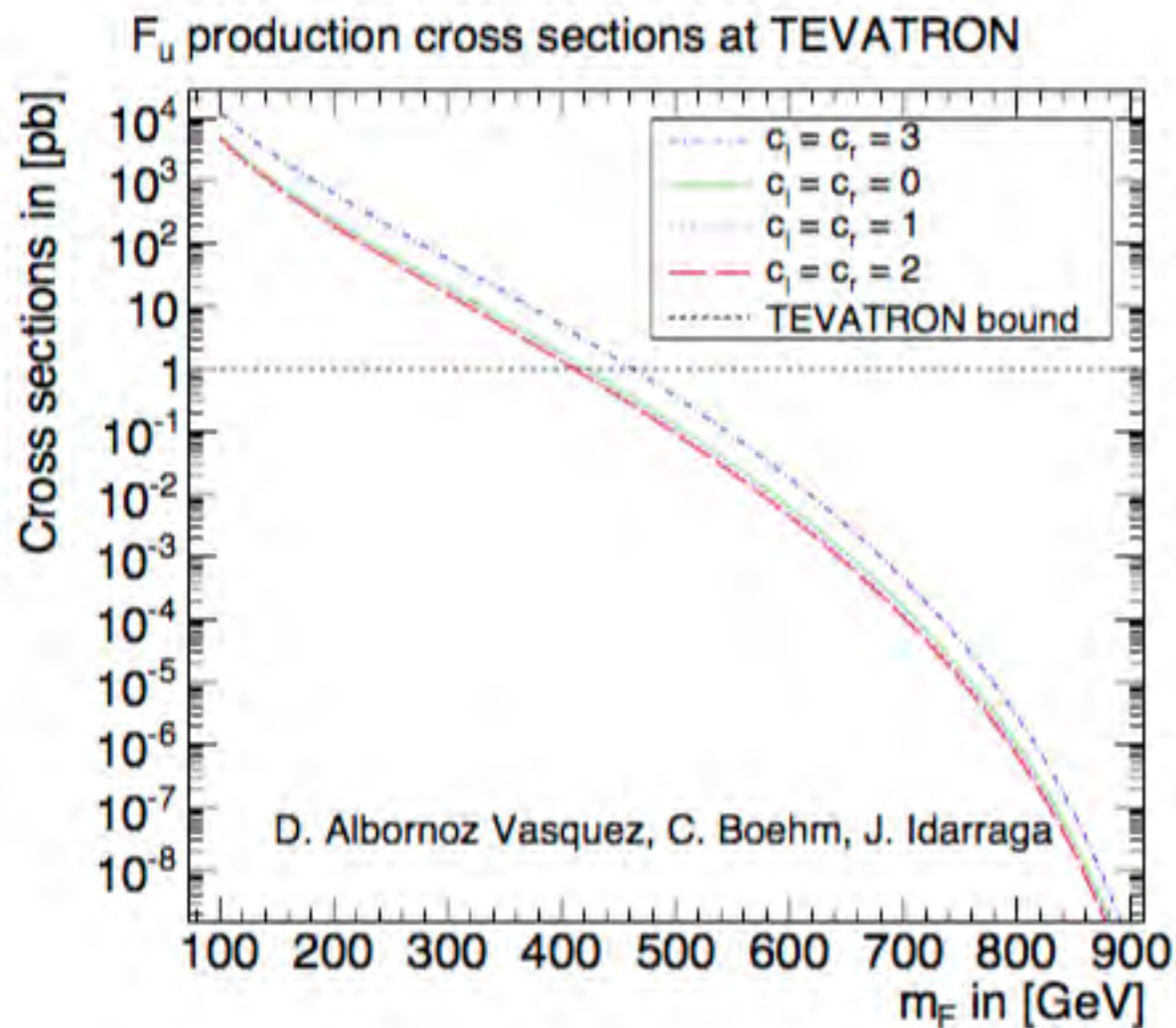
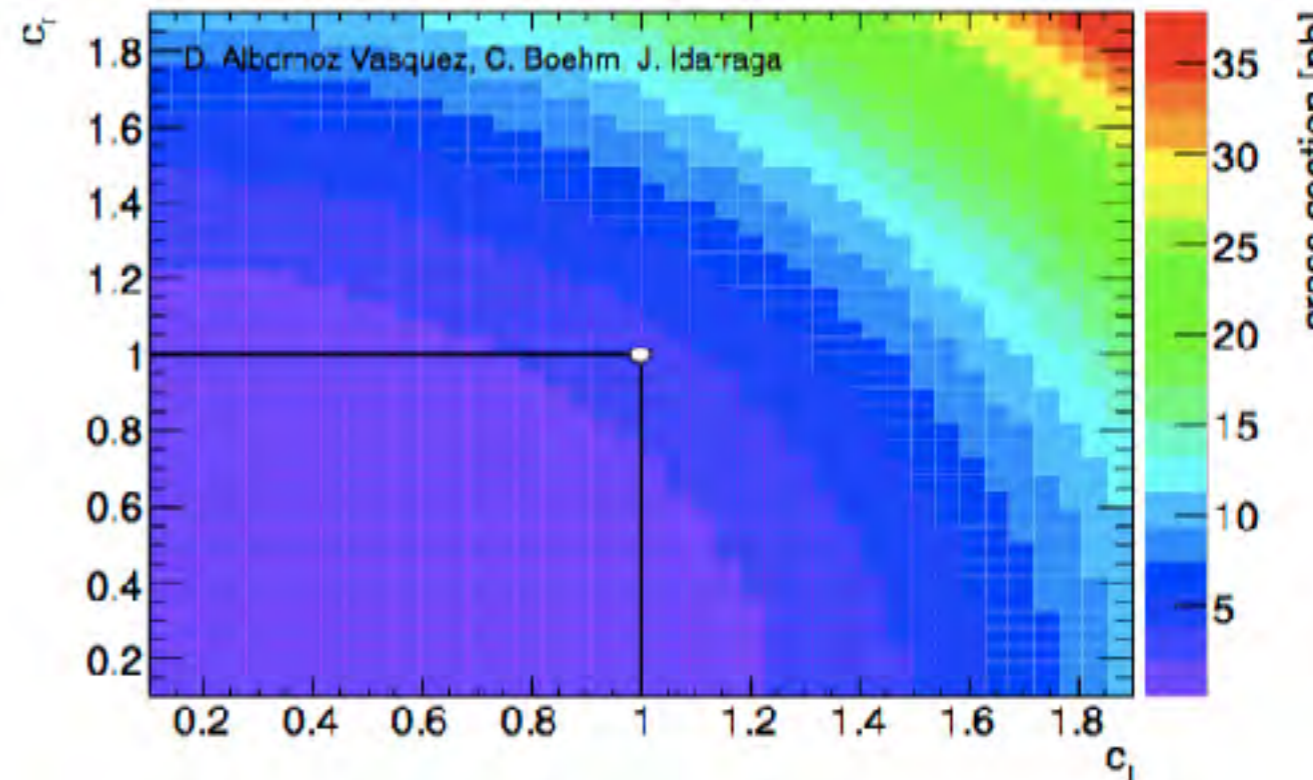
In the early Universe ($c=1$), light DM means light Z' for thermal RD and adjust couplings

In late Universe, light DM is ~ fine because the cross section is velocity dependent

Signature of light DM (t-channel mediators) at LHC



0912.5373

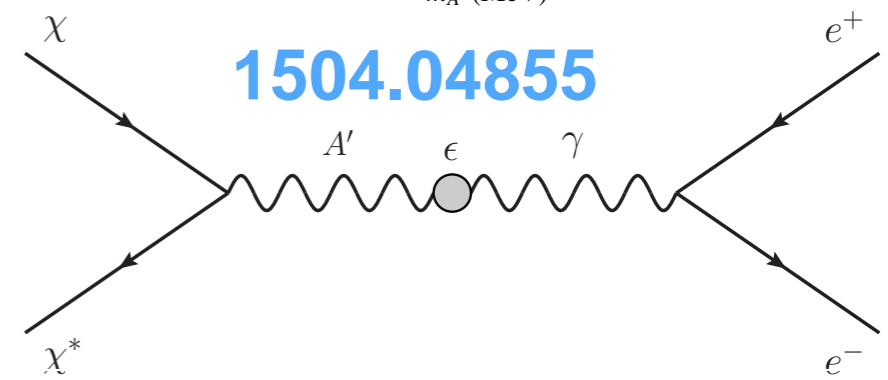
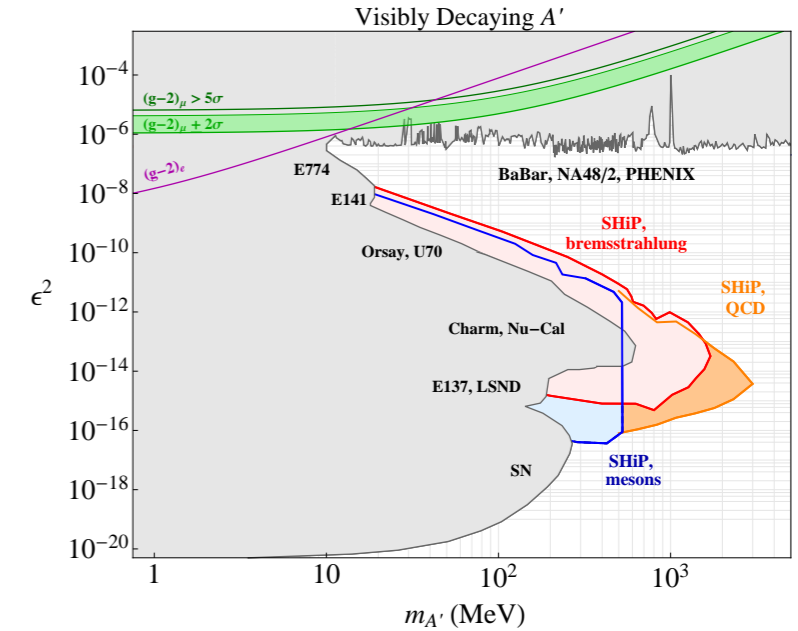
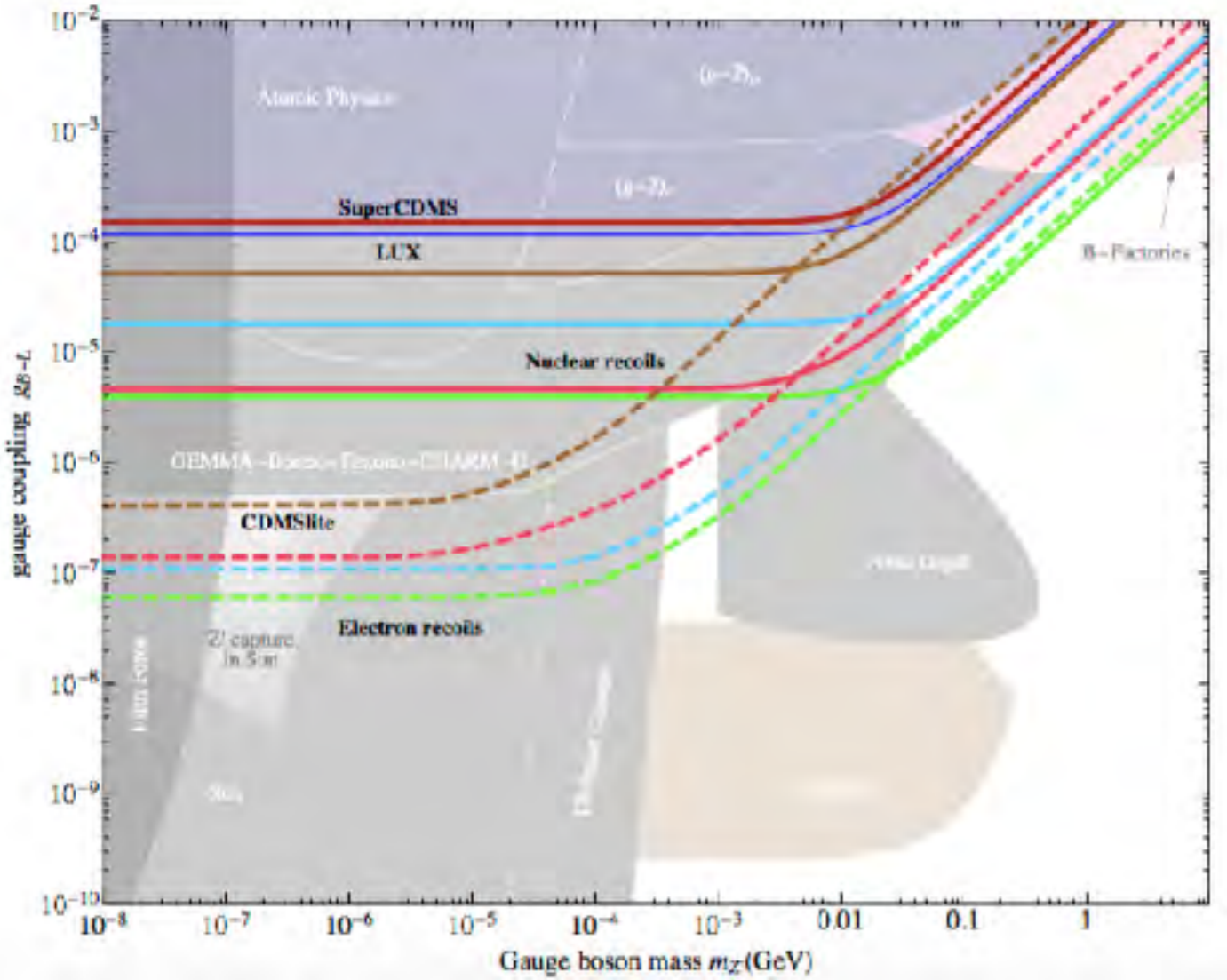


first example of simplified models at LHC

The mediator can be produced through the exchange of DM

Strong constraints!

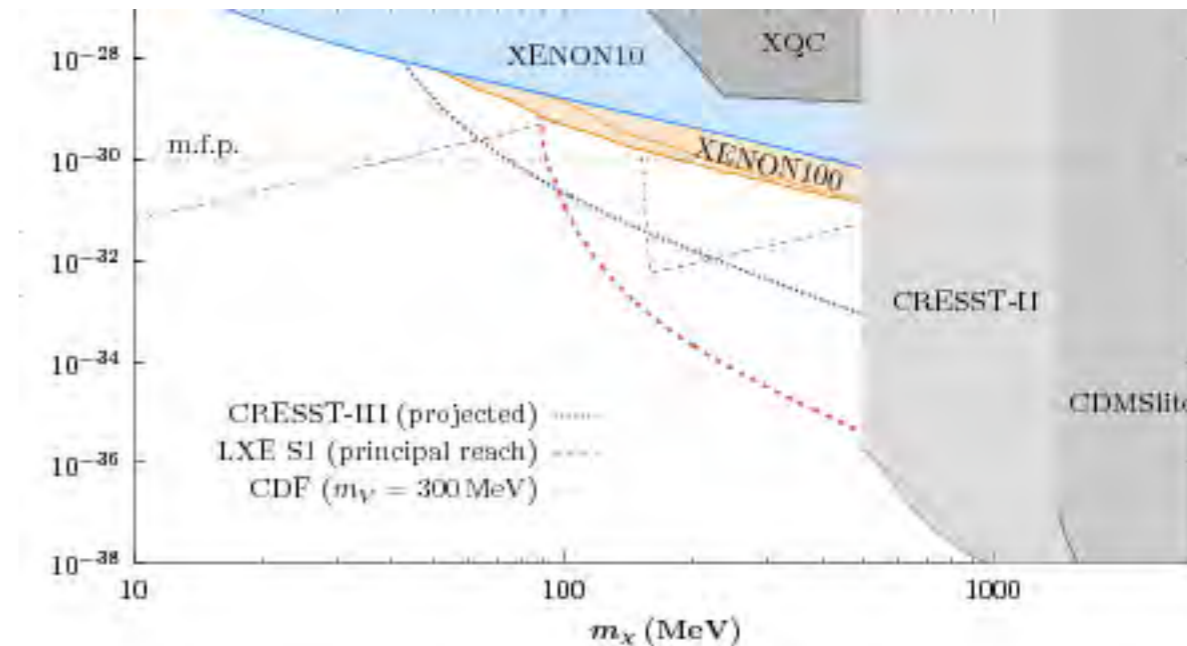
But Light DM particles are being ruled out



1607.01789

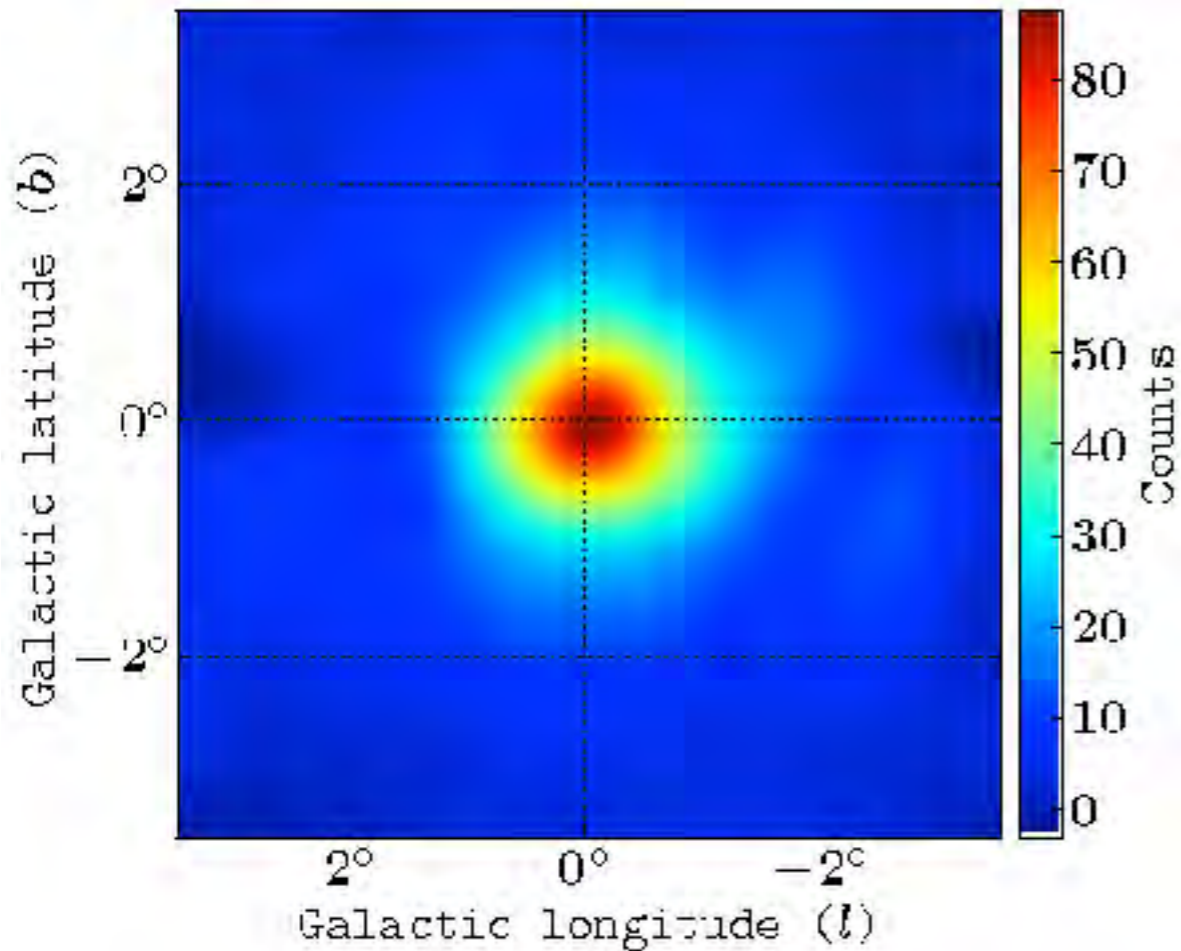
[BDX arXiv:1406.3028](#), [Dafne/Kloe, arXiv:1606.08849](#) & [arXiv:1604.08206](#)

[arXiv:1412.8378](#) [arXiv:1407.0993](#)



The GeV excess (still room for relatively light DM)

FERMI-LAT data

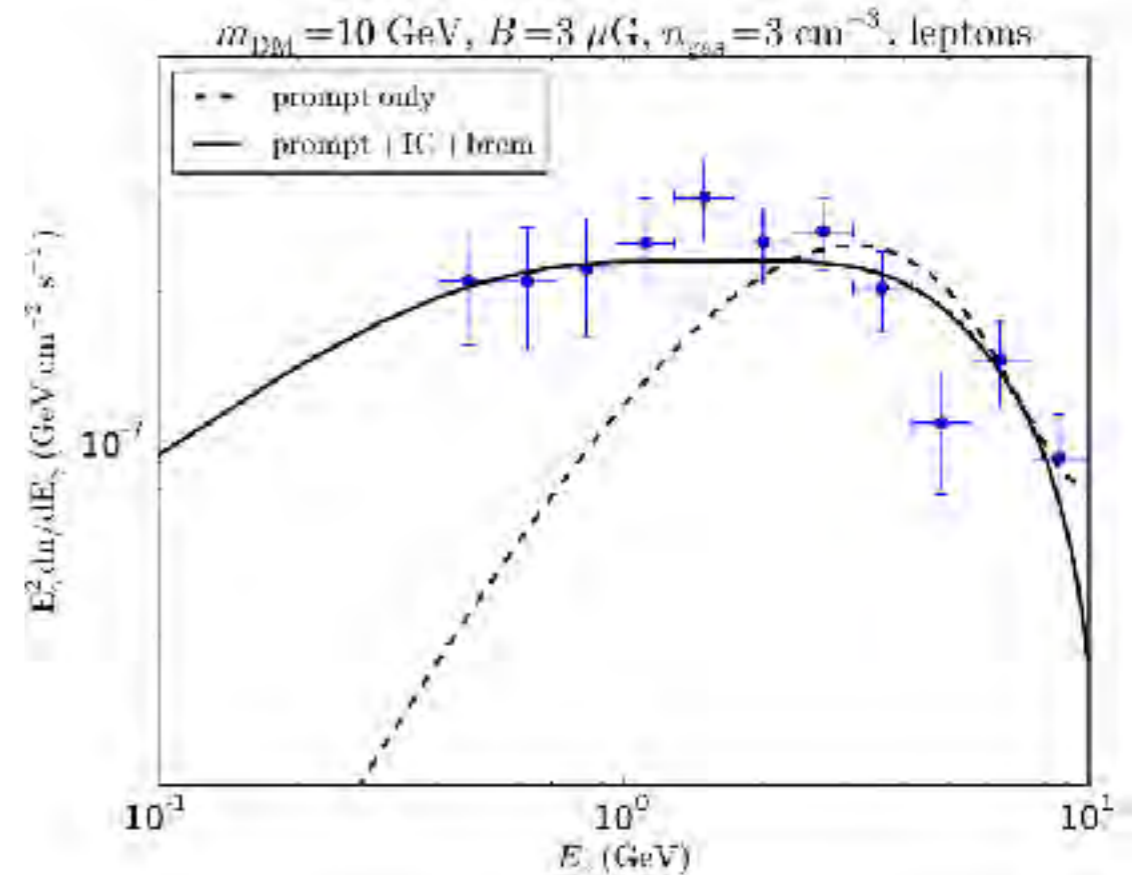


**10-30 GeV DM annihilating mostly
into b-quarks or muons**

Hooper&Goodenough 2009

FERMI-LAT 2009

arXiv:1306.5725

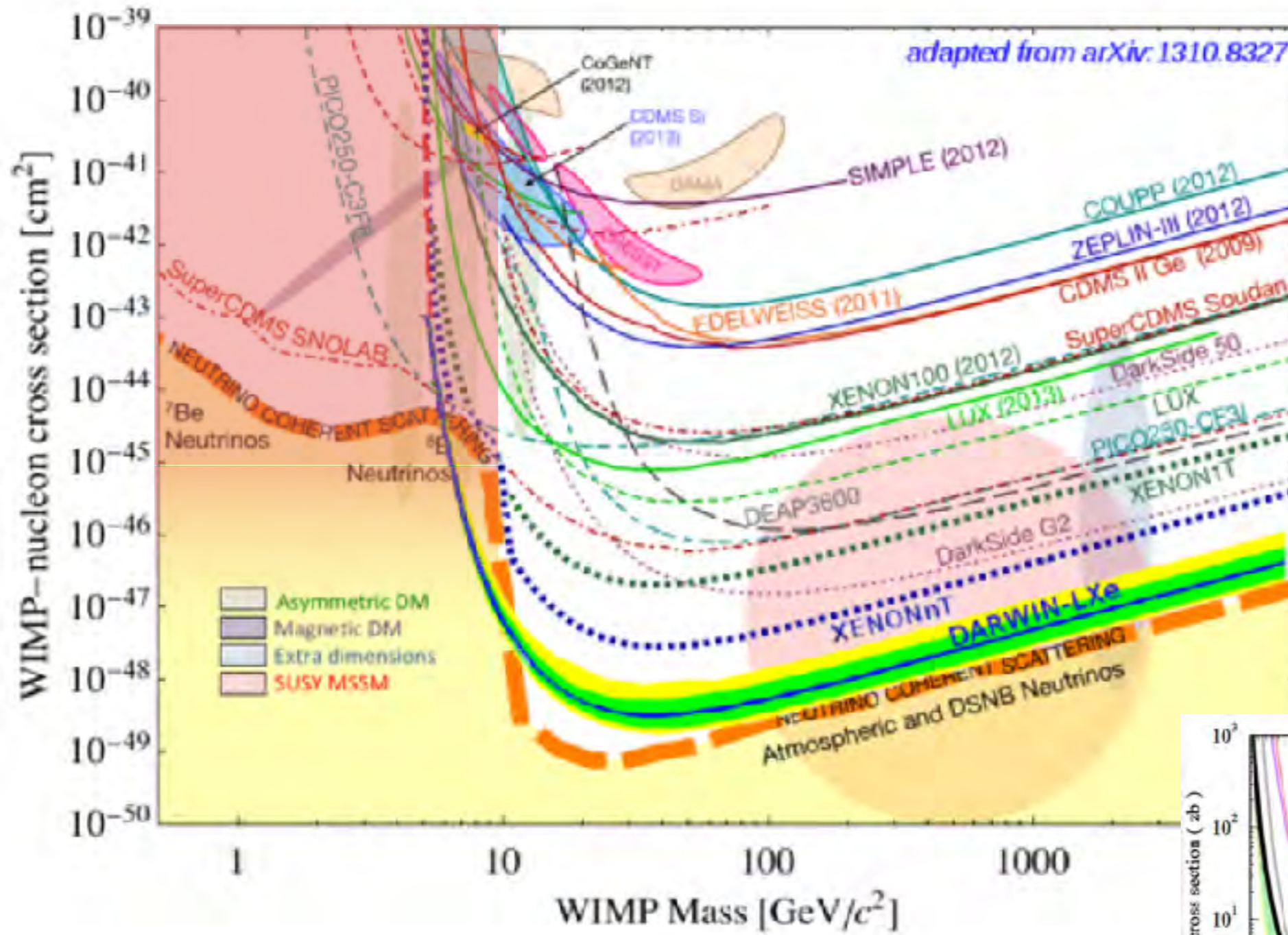


Leptons work too...

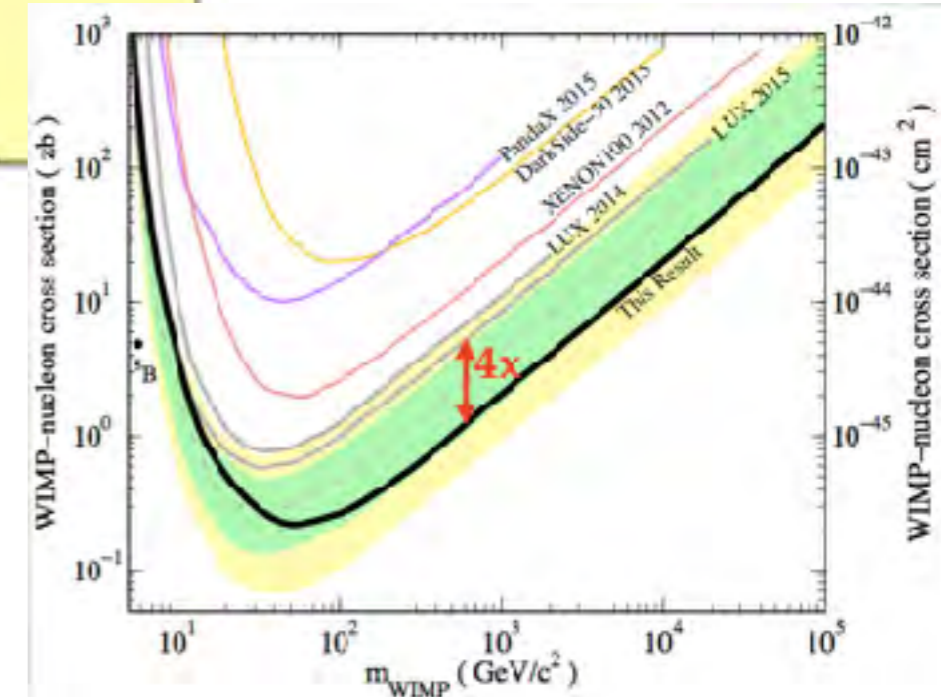
T. Lacroix, CB, J. Silk, 2014

Probably astrophysical sources but ...

Direct detection experiments can be compared to ID and LHC searches!



$$\eta(E, t) = \int_{v_{\min}(E)}^{\infty} \frac{f(v, u_e(t))}{v} d^3v$$



Simplified models

hep-ph/0305261

	Scalar	Fermion	Vector
DM	ϕ	χ	X^μ
Mediator	A	ψ	V^μ
SM (fermions)	—	f	—

Notice the ~!!!

Model Number	DM	Mediator	Interactions	Elastic Scattering	Near Future Reach?	
					Direct	LHC
1	Dirac Fermion	Spin-0	$\bar{\chi}\gamma^5\chi, \bar{f}f$	$\sigma_{SI} \sim (q/2m_\chi)^2$ (scalar)	No	Maybe
1	Majorana Fermion	Spin-0	$\bar{\chi}\gamma^5\chi, \bar{f}f$	$\sigma_{SI} \sim (q/2m_\chi)^2$ (scalar)	No	Maybe
2	Dirac Fermion	Spin-0	$\bar{\chi}\gamma^5\chi, \bar{f}\gamma^5f$	$\sigma_{SD} \sim (q^2/4m_n m_\chi)^2$	Never	Maybe
2	Majorana Fermion	Spin-0	$\bar{\chi}\gamma^5\chi, \bar{f}\gamma^5f$	$\sigma_{SD} \sim (q^2/4m_n m_\chi)^2$	Never	Maybe
3	Dirac Fermion	Spin-1	$\bar{\chi}\gamma^\mu\chi, \bar{b}\gamma_\mu b$	$\sigma_{SI} \sim \text{loop}$ (vector)	Yes	Maybe
4	Dirac Fermion	Spin-1	$\bar{\chi}\gamma^\mu\chi, \bar{f}\gamma_\mu\gamma^5f$	$\sigma_{SD} \sim (q/2m_n)^2$ or $\sigma_{SD} \sim (q/2m_\chi)^2$	Never	Maybe
5	Dirac Fermion	Spin-1	$\bar{\chi}\gamma^\mu\gamma^5\chi, \bar{f}\gamma_\mu\gamma^5f$	$\sigma_{SD} \sim 1$	Yes	Maybe
5	Majorana Fermion	Spin-1	$\bar{\chi}\gamma^\mu\gamma^5\chi, \bar{f}\gamma_\mu\gamma^5f$	$\sigma_{SD} \sim 1$	Yes	Maybe
6	Complex Scalar	Spin-0	$\phi^\dagger\phi, \bar{f}\gamma^5f$	$\sigma_{SD} \sim (q/2m_n)^2$	No	Maybe
6	Real Scalar	Spin-0	$\phi^2, \bar{f}\gamma^5f$	$\sigma_{SD} \sim (q/2m_n)^2$	No	Maybe
6	Complex Vector	Spin-0	$B_\mu^\dagger B^\mu, \bar{f}\gamma^5f$	$\sigma_{SD} \sim (q/2m_n)^2$	No	Maybe
6	Real Vector	Spin-0	$B_\mu B^\mu, \bar{f}\gamma^5f$	$\sigma_{SD} \sim (q/2m_n)^2$	No	Maybe
7	Dirac Fermion	Spin-0 (<i>t</i> -ch.)	$\bar{\chi}(1 \pm \gamma^5)b$	$\sigma_{SI} \sim \text{loop}$ (vector)	Yes	Yes
7	Dirac Fermion	Spin-1 (<i>t</i> -ch.)	$\bar{\chi}\gamma^\mu(1 \pm \gamma^5)b$	$\sigma_{SI} \sim \text{loop}$ (vector)	Yes	Yes
8	Complex Vector	Spin-1/2 (<i>t</i> -ch.)	$X_\mu^\dagger\gamma^\mu(1 \pm \gamma^5)b$	$\sigma_{SI} \sim \text{loop}$ (vector)	Yes	Yes
8	Real Vector	Spin-1/2 (<i>t</i> -ch.)	$X_\mu\gamma^\mu(1 \pm \gamma^5)b$	$\sigma_{SI} \sim \text{loop}$ (vector)	Yes	Yes

arXiv:1404.0022

New types of DM-nuclei interactions

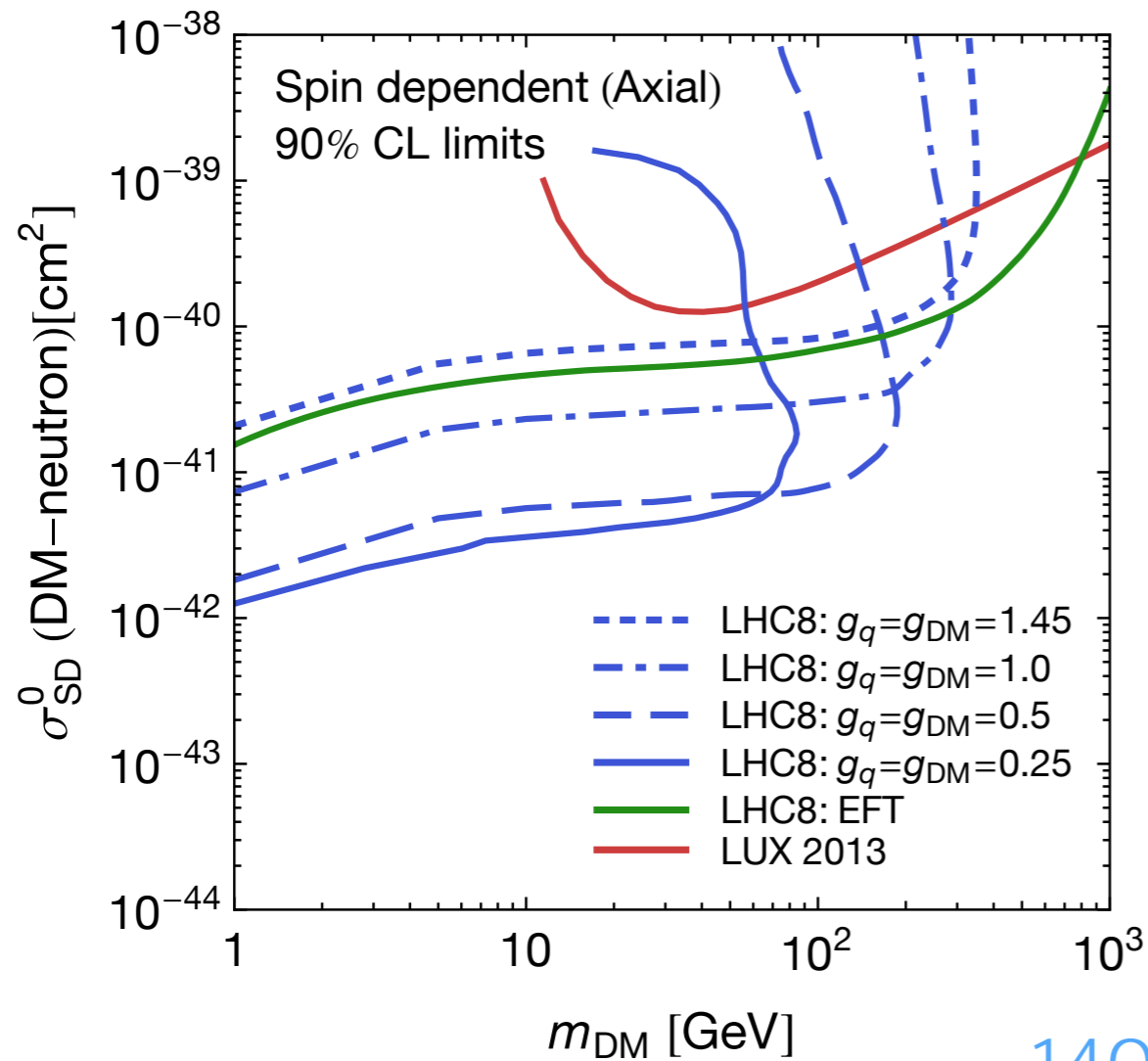
new models, new operators, new couplings to nuclei!

	SI \mathcal{O}_1^{NR}	SD \mathcal{O}_4^{NR}	SD \mathcal{O}_6^{NR}	SD \mathcal{O}_7^{NR}	SI \mathcal{O}_8^{NR}	SD \mathcal{O}_9^{NR}	SD \mathcal{O}_{10}^{NR}	SI \mathcal{O}_{11}^{NR}	SD \mathcal{O}_{13}^{NR}	SD \mathcal{O}_{14}^{NR}
1/2-S	$g_{f,s}g_{\chi,s}$		$g_{f,p}g_{\chi,p}$				$g_{f,p}g_{\chi,s}$	$g_{f,s}g_{\chi,p}$		$g_{f,s}g_{\chi,p}$
1/2*-V	$g_{f,v}g_{\chi,v}$	$g_{f,a}g_{\chi,a}$		$g_{f,a}g_{\chi,v}$	$g_{f,v}g_{\chi,a}$	$g_{f,a}g_{\chi,v}$				
1/2-V		$g_{f,a}g_{\chi,a}$			$g_{f,v}g_{\chi,a}$					
1/2-S $^\pm$	$ g_s^2 - g_p^2 $	$ g_s^2 + g_p^2 $								
1/2*-S $^\pm$	$ g_s^2 \pm g_p^2 $	$ g_s^2 + g_p^2 $								
1/2-V $^\pm$	$ g_v^2 - g_a^2 $	$ g_v^2 + g_a^2 $								
1/2*-V $^\pm$	$ a_v^2 \pm a_a^2 $	$ a_v^2 + a_a^2 $								
0-S	$g_\chi g_{f,s}$						$g_\chi g_{f,p}$			
0-V	$g_\chi g_{f,v}$			$g_\chi g_{f,a}$			$g_\chi g_{f,u}$			
0-F $^\pm$	$ g_s ^2 \pm g_p ^2$						$g_p g_s^\dagger - g_s g_p^\dagger$			
0-F $^\pm$	$ g_s ^2 \pm g_p ^2$						$g_p g_s^\dagger \pm g_s g_p^\dagger$			
1-S	$g_\chi g_{f,s}$						$g_\chi g_{f,p}$			
1-V (\mathcal{V}_1)		$\text{Im}(g_\chi)g_v$	$\text{Im}(g_\chi)g_v$		$\text{Im}(g_\chi)g_a$			$\text{Re}(g_\chi)g_a$	$\text{Re}(g_\chi)g_v$	
1-V (\mathcal{V}_2)	$\text{Im}(g_\chi)g_v$			$\text{Im}(g_\chi)g_a$			$\text{Re}(g_\chi)g_a$			
1-V (\mathcal{V}_3)		$\text{Re}(g_\chi)g_a$	$\text{Re}(g_\chi)g_a$		$\text{Re}(g_\chi)g_v$	$\text{Re}(g_\chi)g_v$		$\text{Im}(g_\chi)g_v$		$\text{Im}(g_\chi)g_a$
1-F $^\pm$	$ g_s ^2 \pm g_p ^2$				$g_v g_a^\dagger + g_a g_v^\dagger$	$g_v g_a^\dagger + g_a g_v^\dagger$	$g_v g_a^\dagger - g_a g_v^\dagger$			
1*-F $^\pm$	$ g_s ^2 \pm g_p ^2$					$g_v g_a^\dagger + g_a g_v^\dagger$	$g_v g_a^\dagger - g_a g_v^\dagger$	$g_v g_a^\dagger + g_a g_v^\dagger$		

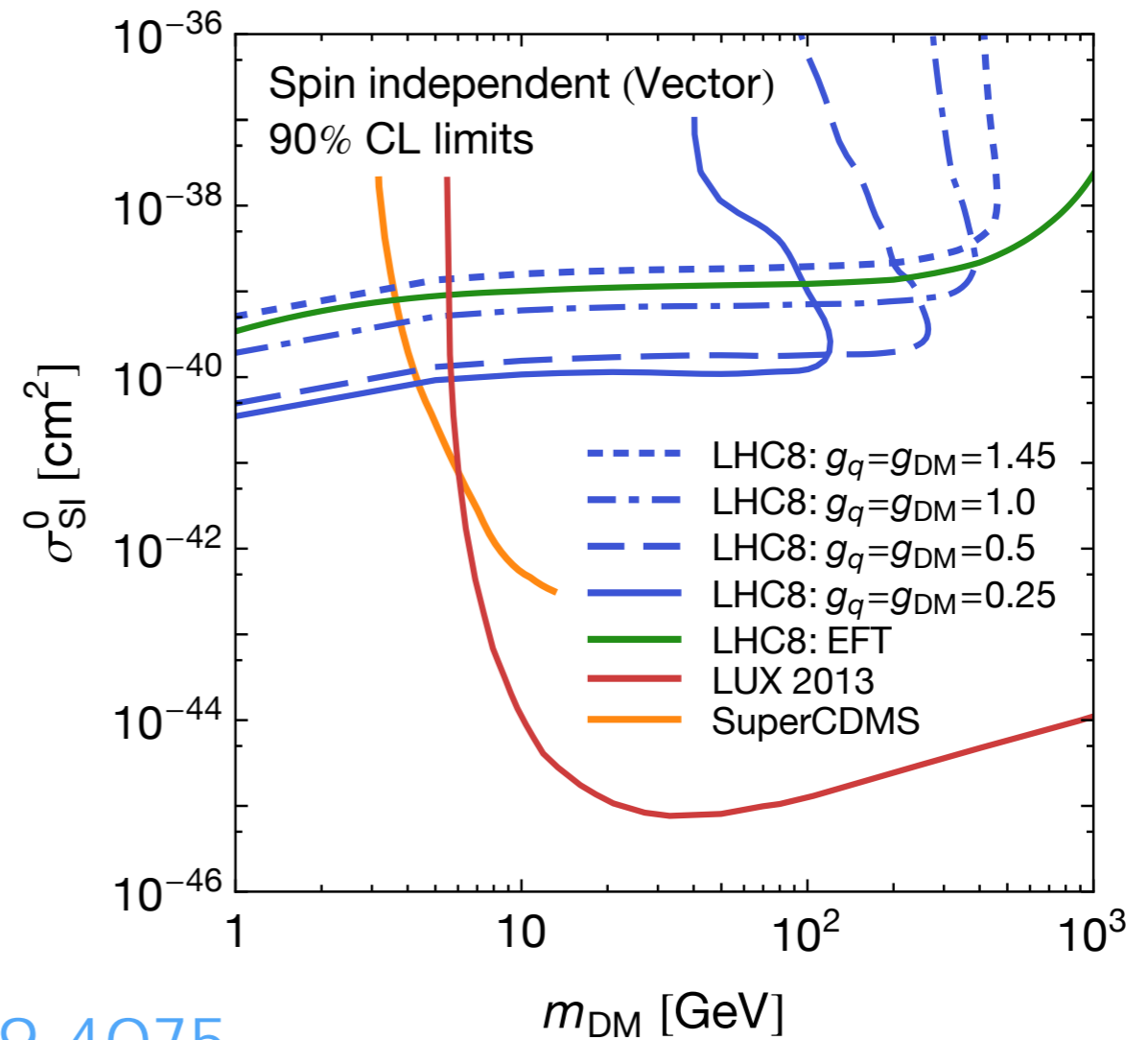
main operators

new operators

most of them suppressed by spin and/or velocity but ..



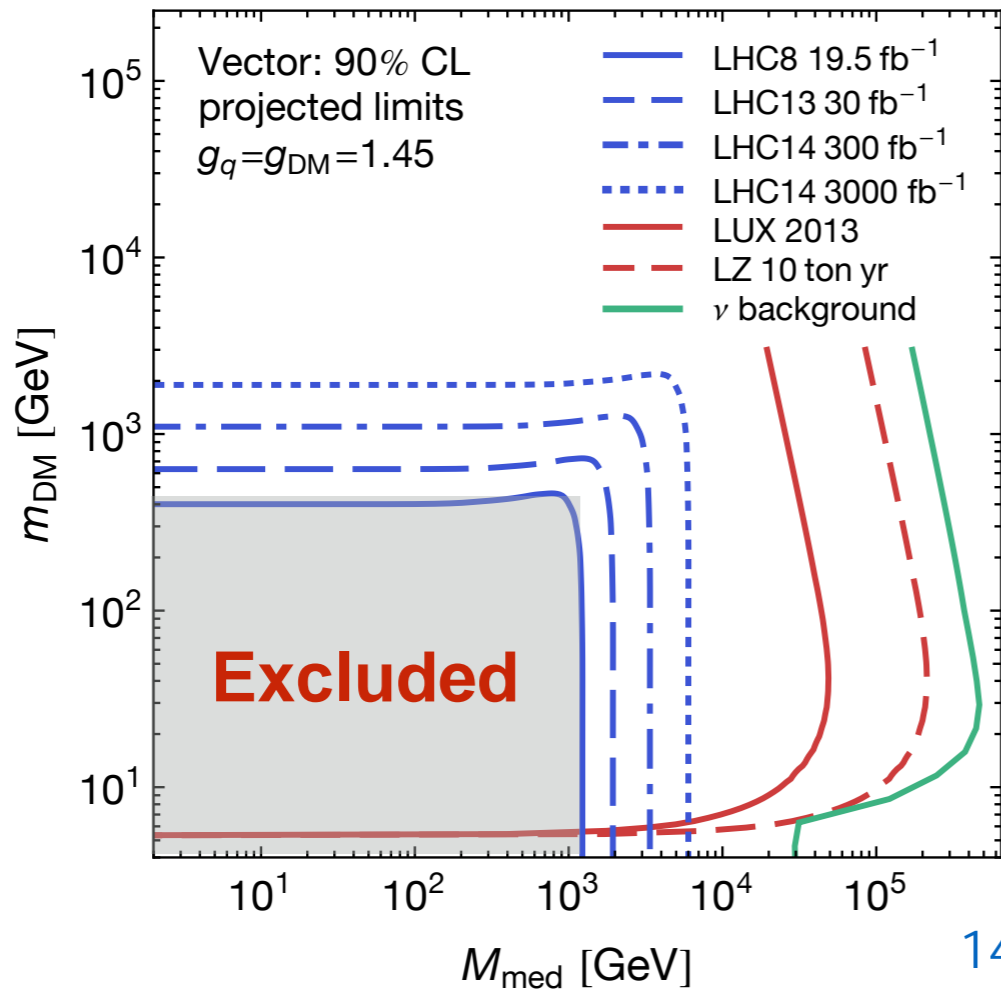
1409.4075



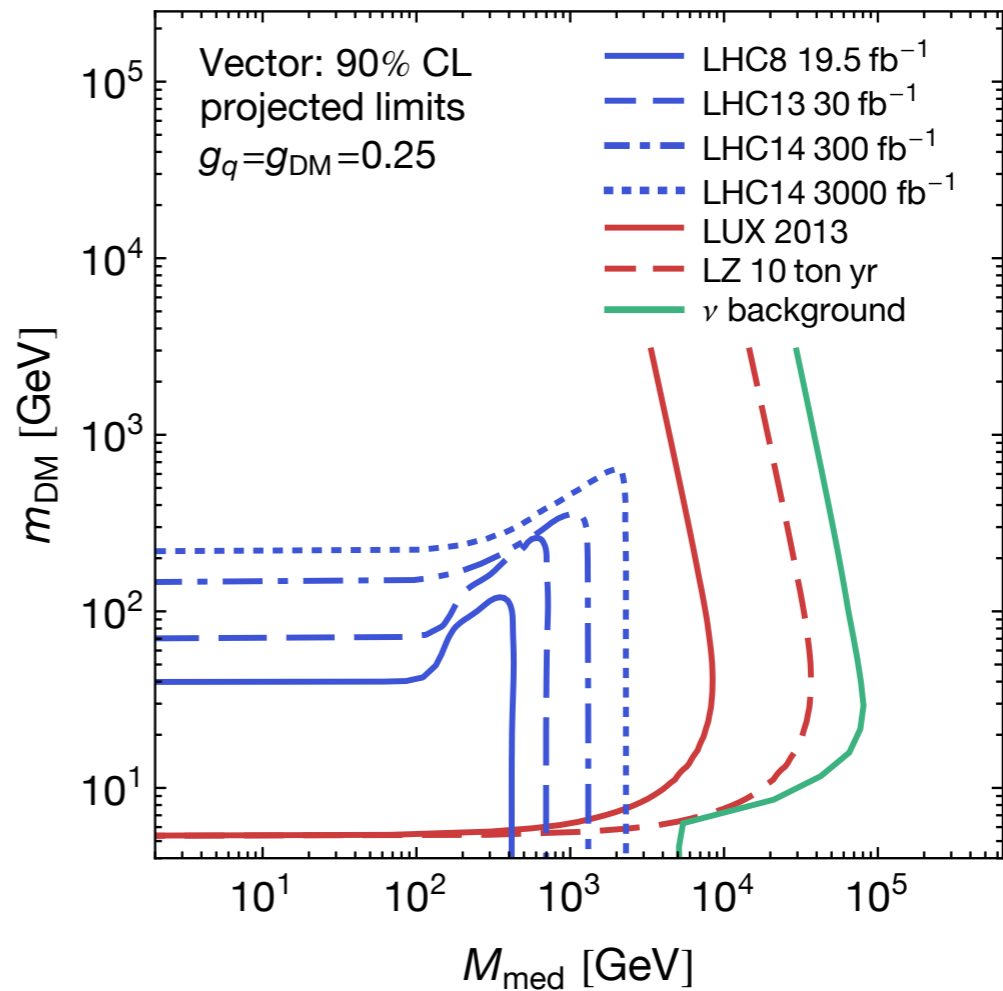
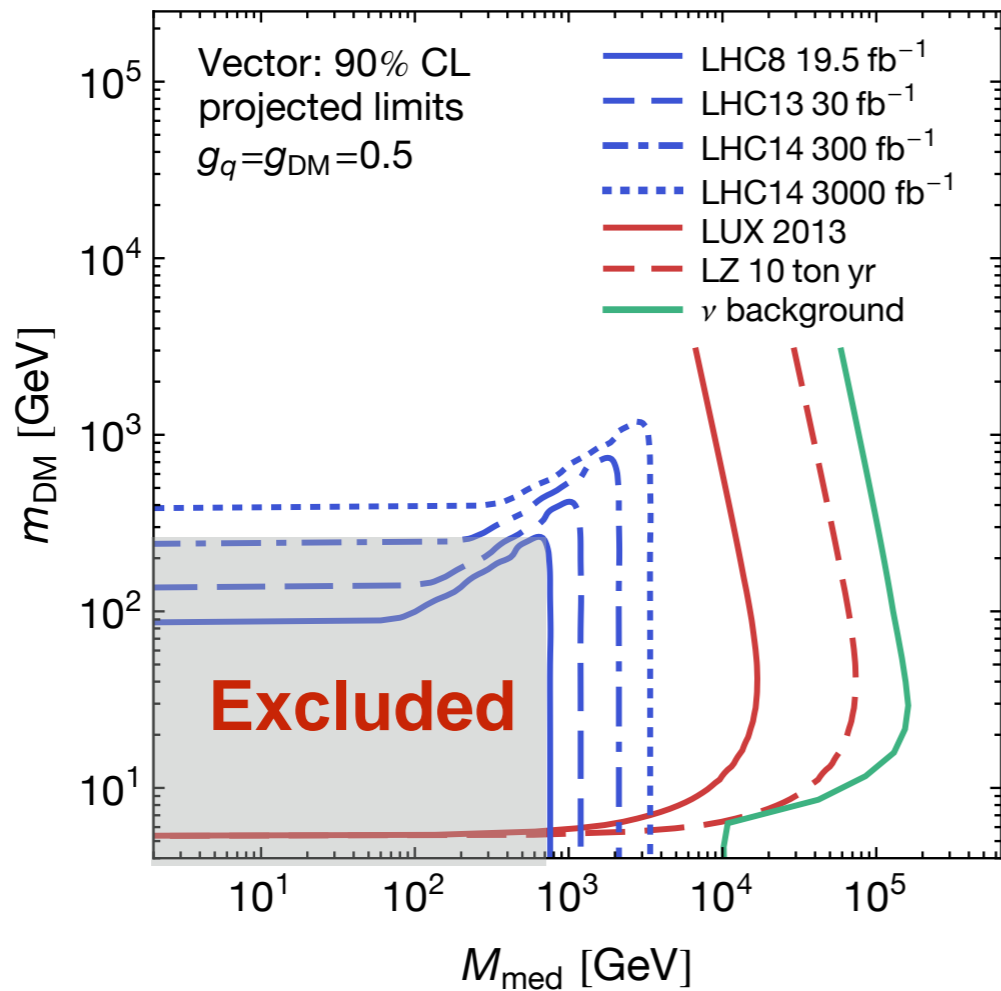
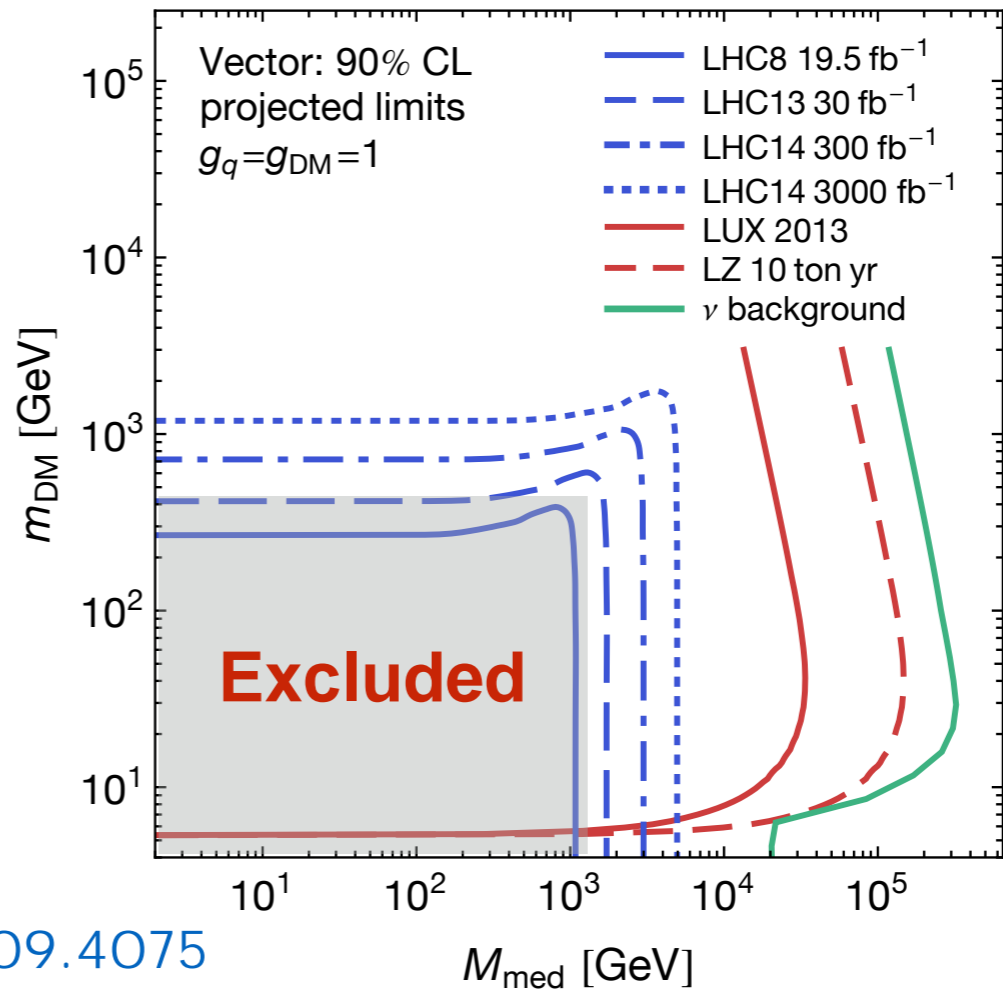
Competition between DD and LHC (monojets)

EFT not valid at low mass so simplified models!

no resonance



1409.4075



Conclusion

Stringent limits from Direct and Indirect detection

Still room for new physics but constrained!

We need to pay attention to the cosmology

as this can affect all our estimates!